

CP Violation in the B_s^0 System at the Tevatron



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On Behalf of the CDF and DØ Collaborations



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Introduction to CPV in B_s^0 Decays

CPV studies in B_s^0 decays aim to understand the source of CP Violation.

One of the last remaining places to search for New Physics!

The success of the B -factories has shown that large ($> \sim 10\%$) contributions of NP are excluded from tree-level B^+ and B^0 decays.

B_s^0 decays are much less constrained and even if the B_s^0 mixing measurement constrains the strength of NP, current experimental knowledge does not exclude large (up to π) phases from NP.

CP Violation in B_s^0 is an excellent place to search for NP since it is predicted to be small in the SM. A measurement of a large CP phase is a clear indication of NP.

The Tevatron at Fermilab

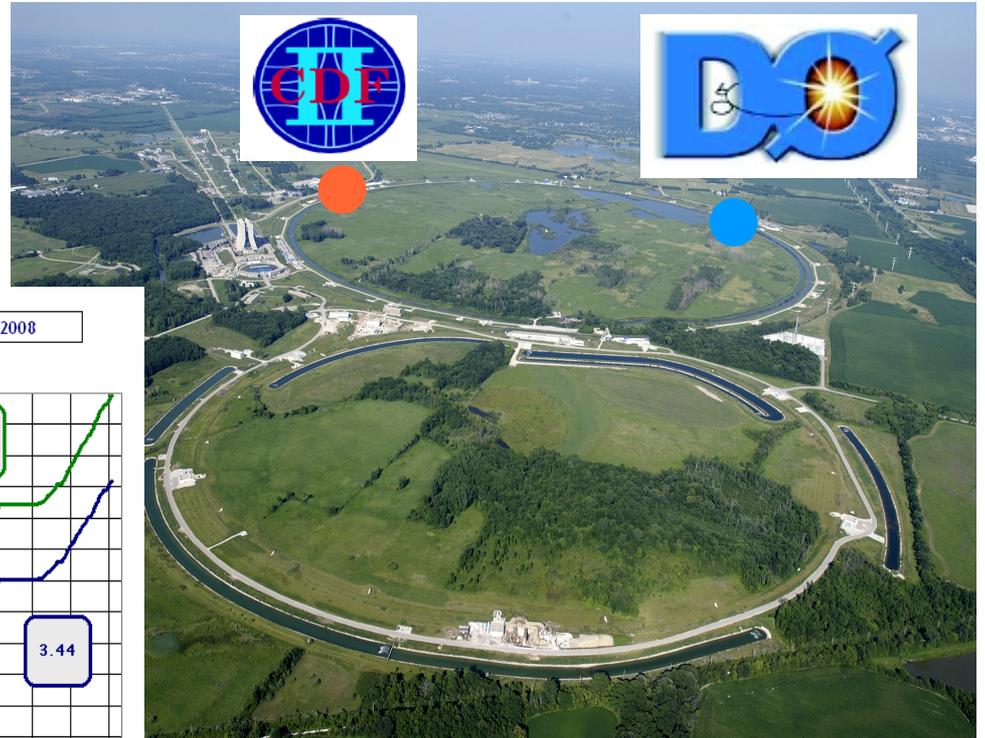
Proton on antiproton collisions at 1.96 TeV energy

Collider Experiments

CDF and **DØ**

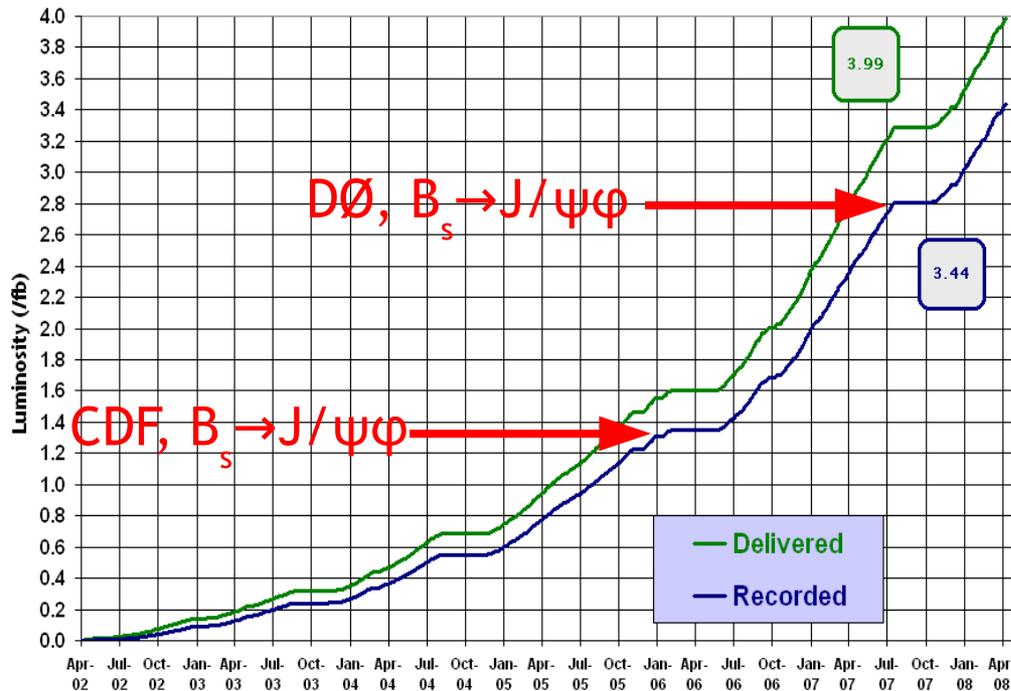
Tevatron running with peak luminosity
 $\sim 315 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$

$\sim 3.9 \text{ fb}^{-1}$ delivered, $\sim 3.4 \text{ fb}^{-1}$ recorded
per experiment



Run II Integrated Luminosity

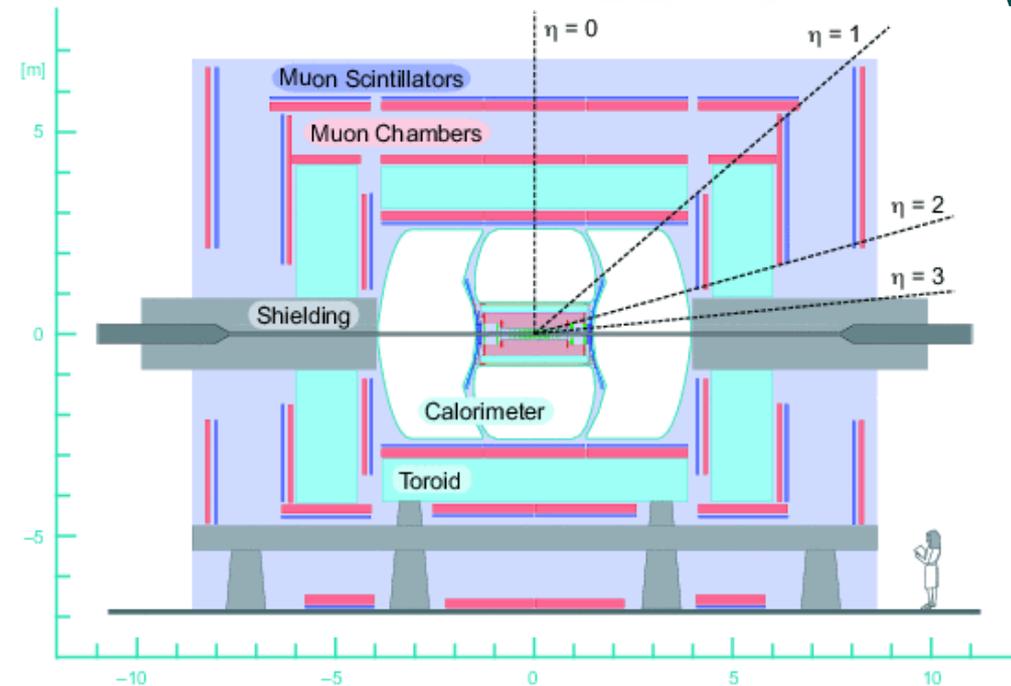
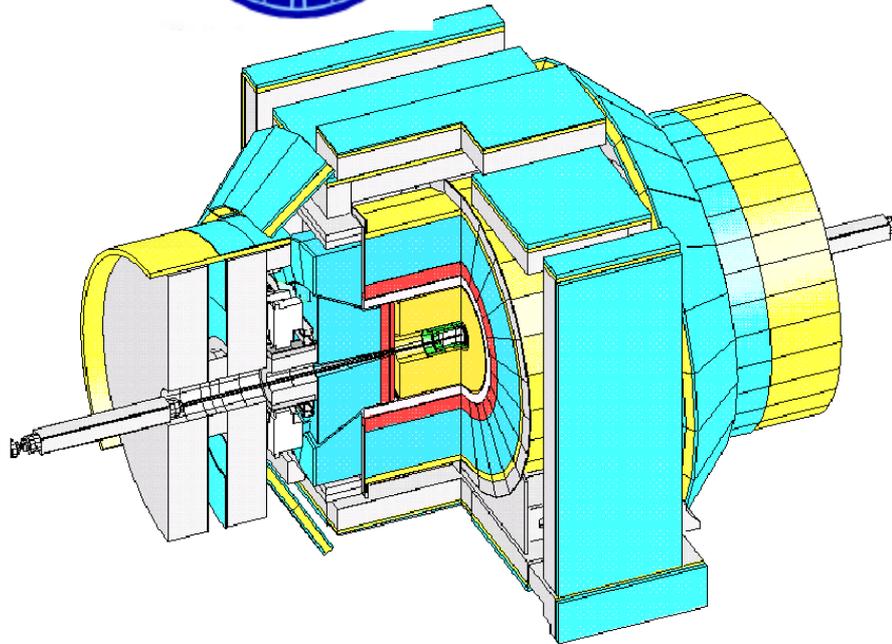
19 April 2002 - 27 April 2008



A lot of data already
collected and waiting
to be analyzed!



Collider Detectors



Relevant for B physics

CDF Tracker: mass resolution, vertexing

Silicon & L00

Large radii drift chamber

excellent momentum resolution

dE/dx and particle id

Triggered Muon Coverage: $|\eta| < 1$

Time of flight, particle ID

DØ Tracker: coverage, vertexing

Silicon & scintillating fiber

Small radii, $|\eta| < 2$

New layer 0 silicon on beam pipe

Improves impact parameter res.

Triggered Muon Coverage: $|\eta| < 2$

Single muon

Di-muon

$B_s^0 - \bar{B}_s^0$ Mixing

Flavor eigenstates propagate according to the Schrodinger Eq.

$$i \frac{d}{dt} \begin{pmatrix} B_s(t) \\ \bar{B}_s(t) \end{pmatrix} = \left(\begin{bmatrix} m & M_{12}^s \\ M_{12}^{s*} & m \end{bmatrix} - \frac{i}{2} \begin{bmatrix} \Gamma & \Gamma_{12}^s \\ \Gamma_{12}^{s*} & \Gamma \end{bmatrix} \right) \begin{pmatrix} B_s(t) \\ \bar{B}_s(t) \end{pmatrix}$$

Diagonalizing gives two physically observed
“*Heavy*” and “*Light*” mass eigenstates

$$|B_s^H\rangle = p |B_s^0\rangle - q |\bar{B}_s^0\rangle \quad |B_s^L\rangle = p |B_s^0\rangle + q |\bar{B}_s^0\rangle$$

Observables

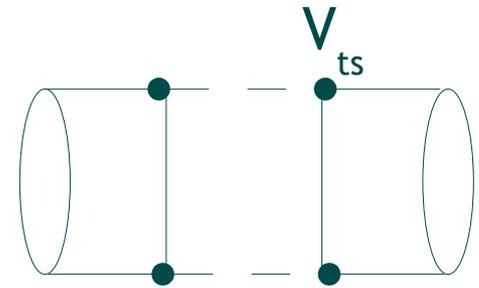
$$\Delta M_s = M_H - M_L \approx 2|M_{12}|$$

$$\Delta \Gamma^{\text{CP}} = \Gamma_{\text{even}} - \Gamma_{\text{odd}} \approx 2|\Gamma_{12}|$$

$$\Delta \Gamma_s = \Gamma_L - \Gamma_H \approx 2|\Gamma_{12}| \cos(\varphi_s) \quad \varphi_s = \arg(-M_{12}/\Gamma_{12})$$

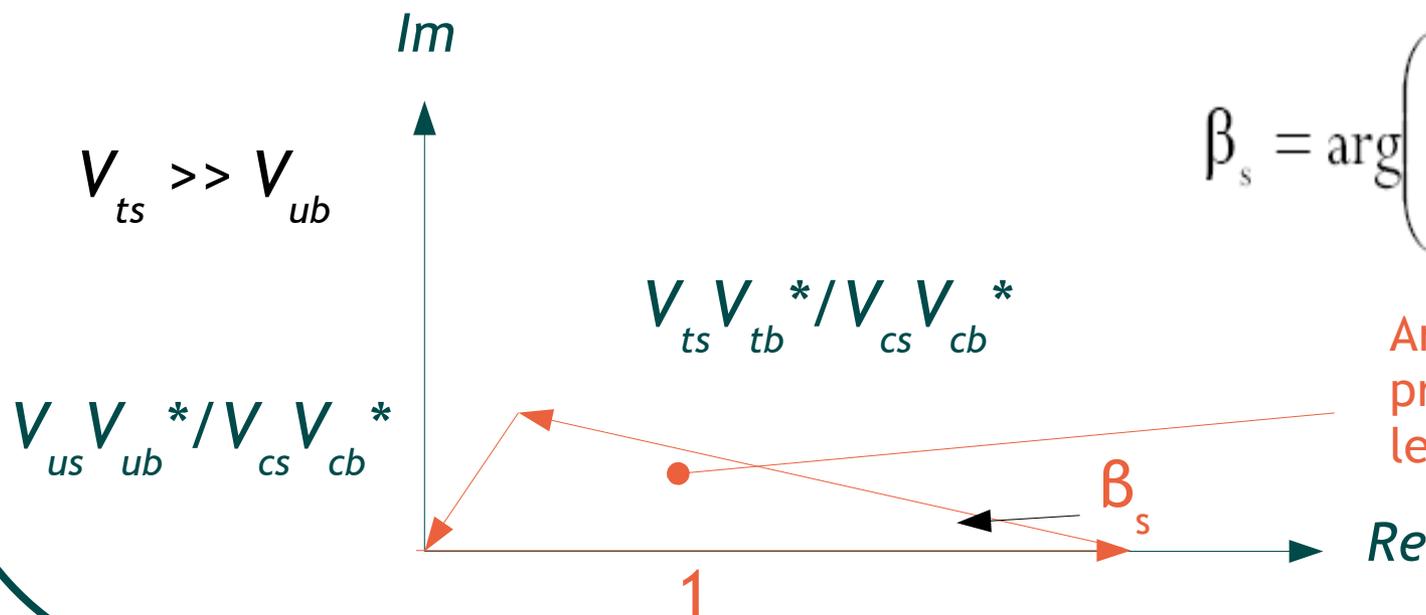
CP Violation in the B_s^0 System

CKM Matrix $\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$



SM accommodates CPV by introducing a single complex phase in the CKM matrix

B_s^0 unitary condition $V_{us}V_{ub}^* + V_{cs}V_{cb}^* + V_{ts}V_{tb}^* = 0$

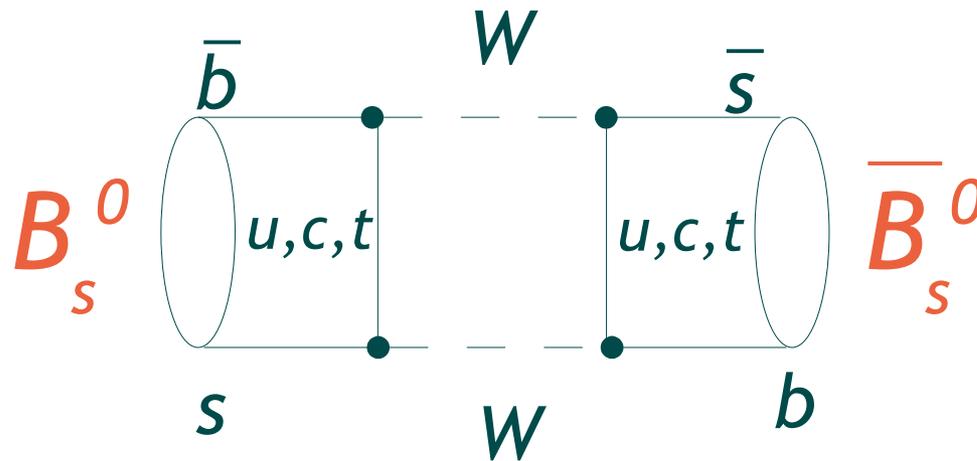


$$\beta_s = \arg\left(-\frac{V_{ts}V_{tb}^*}{V_{cs}V_{cb}^*}\right) \sim 0.02$$

Area of triangle proportional to level of CP violation

CP Violation in $B_s^0 - \bar{B}_s^0$ Mixing

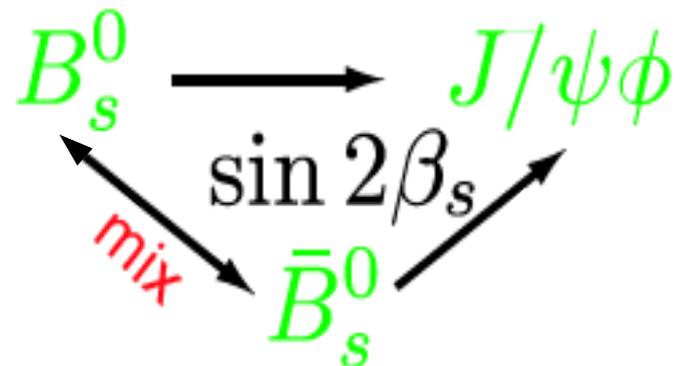
Matter \longleftrightarrow Antimatter



Semileptonic asymmetry

$$N(B_s^0 \rightarrow D^{*-}) \text{ vs } N(\bar{B}_s^0 \rightarrow B_s^0 \rightarrow D^{*-})$$

Interference



CP Violation in the B_s^0 System

How could new physics affect these phases?

$$2\beta_s^{SM} \rightarrow \overset{\sim 0.04}{2\beta_s^{SM}} - \phi_s^{NP}$$

$$\phi_s^{SM} = \arg[-M_{12}/\Gamma_{12}] \rightarrow \underset{\sim 0.004}{\phi_s^{SM}} + \phi_s^{NP}$$

Both CDF and $D\bar{0}$ measure the phase responsible for CP violation in $B_s^0 \rightarrow J/\psi\phi$ decays

$$\underset{D\bar{0}}{\phi_s^{J/\psi\phi}} = \underset{CDF}{-2\beta_s^{J/\psi\phi}} \approx \underset{\text{If large}}{\phi_s^{NP}}$$

Topics Covered in this Talk

Semileptonic Asymmetry

DØ, 1.3 fb⁻¹ PRL 98, 151801 (2007)

DØ, 1.0 fb⁻¹ PRD 74, 092001 (2006)

Combined: PRD 76, 057101 (2007)

CDF, 1.6 fb⁻¹ CDF note 9015

Mixing and Decay $B_s^0 \rightarrow J/\psi\phi$

DØ, 1.1 fb⁻¹ PRL 98, 121801 (2007)

DØ, 2.8 fb⁻¹ submitted to PRL

CDF, 1.35 fb⁻¹ PRL 100, 161802 (2008)

CP Violation: Semileptonic Asymmetry

Measured semileptonic asymmetries

DØ, 1.3 fb⁻¹, PRL 98, 151801 (2007)

$$N(B_s^0 \rightarrow D_s^- \mu^+ \nu)$$

vs.

$$N(\bar{B}_s^0 \rightarrow D_s^+ \mu^- \bar{\nu})$$

+

DØ, 1 fb⁻¹, PRD 74, 092001 (2006)

$$N(b\bar{b} \rightarrow \mu^+ \mu^+ X)$$

vs.

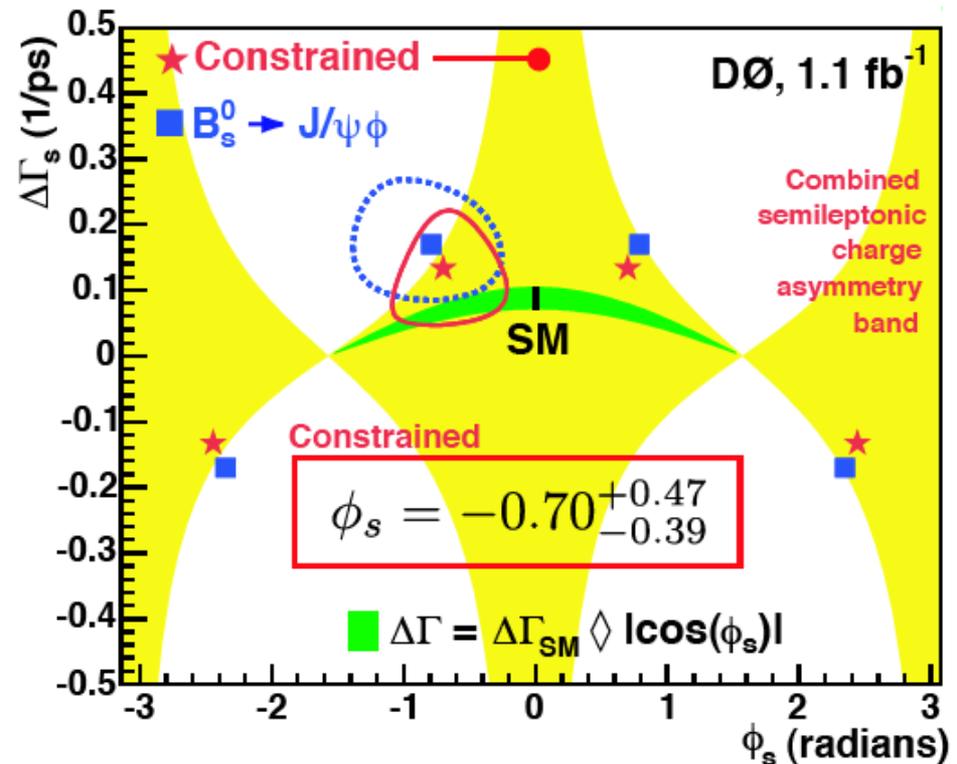
$$N(b\bar{b} \rightarrow \mu^- \mu^- X)$$

||

$$a_{SL}^s = 0.0001 \pm 0.0090$$

PRD 76, 057101 (2007)

$$a_{sl}^s \equiv \frac{N(\bar{B}_s \rightarrow f) - N(B_s \rightarrow \bar{f})}{N(\bar{B}_s \rightarrow f) + N(B_s \rightarrow \bar{f})} = \frac{\Delta\Gamma_s}{\Delta M_s} \tan \phi_s$$



Regular flipping of polarity of solenoid (tracking)
and toroid (muons) helps control systematic uncertainties.

CP Violation: Semileptonic Asymmetry

Measured semileptonic asymmetries

$D\bar{D}$, 1.3 fb^{-1} , PRL 98, 151801 (2007)

$$N(B_s^0 \rightarrow D_s^- \mu^+ \nu) \quad a_{sl}^s \equiv \frac{N(\bar{B}_s \rightarrow f) - N(B_s \rightarrow \bar{f})}{N(\bar{B}_s \rightarrow f) + N(B_s \rightarrow \bar{f})} = \frac{\Delta\Gamma_s}{\Delta M_s} \tan \phi_s$$

vs.

$$N(\bar{B}_s^0 \rightarrow D_s^+ \mu^- \bar{\nu})$$

+

$D\bar{D}$, 1 fb^{-1} , PRD 74, 092001 (2006)

$$N(b\bar{b} \rightarrow \mu^+ \mu^+ X)$$

vs.

$$N(b\bar{b} \rightarrow \mu^- \mu^- X)$$

||

$$a_{SL}^s = 0.0001 \pm 0.0090$$

PRD 76, 057101 (2007)

CDF, 1.6 fb^{-1} , CDF Note 9015

$$N(b\bar{b} \rightarrow \mu^+ \mu^+ X)$$

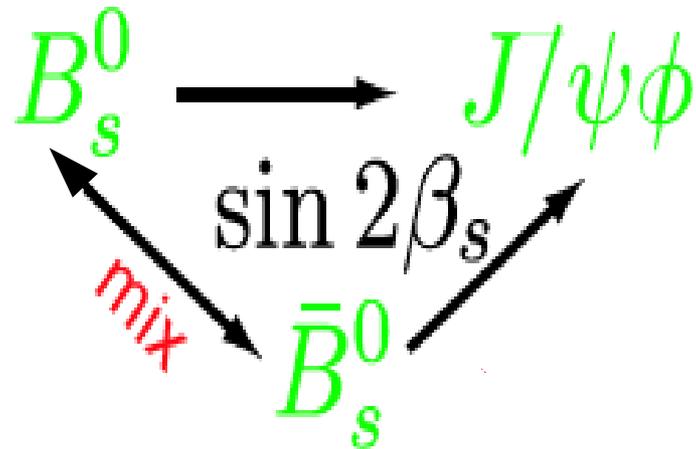
vs.

$$N(b\bar{b} \rightarrow \mu^- \mu^- X)$$

$$a_{SL}^s = 0.020 \pm 0.021 \pm 0.018$$

Regular flipping of polarity of solenoid (tracking)
and toroid (muons) helps control systematic uncertainties.

CP Violation in $B_s^0 \rightarrow J/\psi\Phi$ decays



CP violation becomes observable in these decays due to the interference between the mixing and decay amplitudes.

$J/\psi + \Phi$ is an admixture of states that are both **CP(even)** and **CP(odd)**

Angular analysis is used to separate the CP components and measure the lifetimes of each component

Flavor Tagging gives us useful information on the flavor of the produced B_s^0 meson

$$B_s^0 \rightarrow J/\psi \Phi$$

J/ψ and φ are vector particles and have definite angular distributions for CP-even and CP-odd final states.

$$B_s \rightarrow V1 + V2 (J/\psi + \varphi) \quad \text{Spin } 0 \rightarrow 1 + 1 \quad \ell = 0,1,2$$

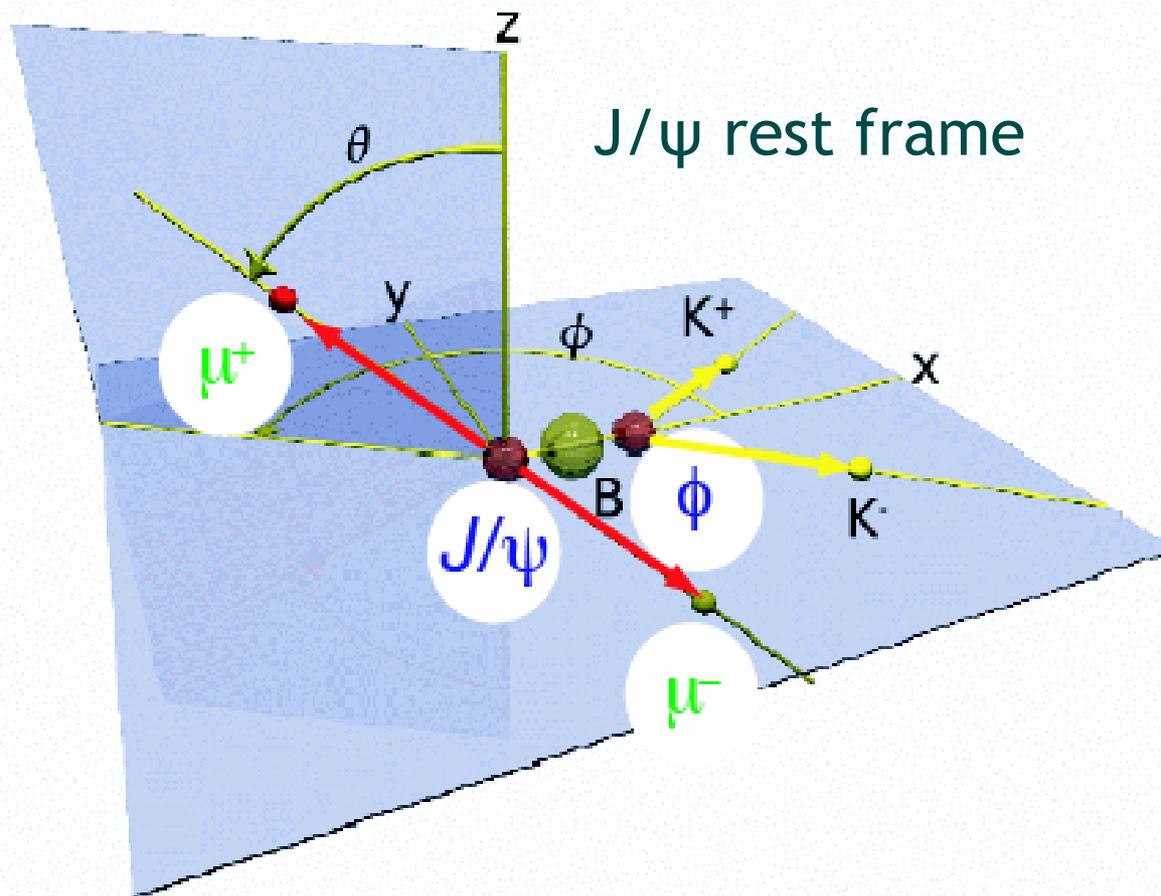
Parameterized angular decay in the Transversity basis.

Angular dependencies are described in terms of polarization amplitudes:

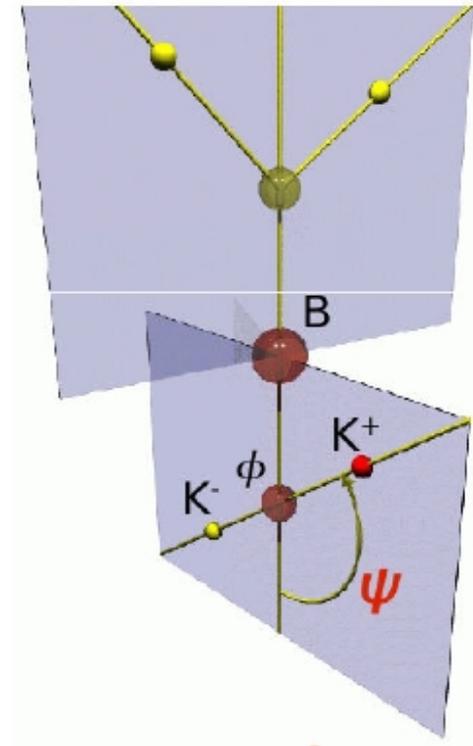
- A_0 : Both vectors longitudinally polarized ($\ell = 0,2$) CP even
- A_{\parallel} : Transversely polarized and vectors parallel ($\ell = 0,2$) CP even
- A_{\perp} : Transversely polarized and vectors perpendicular ($\ell = 1$) CP odd

$$|A_{\parallel}(0)|^2 + |A_{\perp}(0)|^2 + |A_0(0)|^2 = 1$$

Angular Analysis



ϕ rest frame



Angles θ (transversity), ϕ and ψ . ψ is the angle between \vec{p}_{K^+} and the x -axis in the rest frame of ϕ .

Differential Decay Rate and Amplitudes

$$\begin{aligned}
 \frac{d^4\Gamma \left[B_s^0(t) \rightarrow J/\psi(\rightarrow \mu^+\mu^-)\phi(\rightarrow K^+K^-) \right]}{d\cos\theta \, d\varphi \, d\cos\psi \, dt} &\propto \\
 &2\cos^2\psi(1 - \sin^2\theta\cos^2\varphi) \cdot |A_0(t)|^2 \\
 &+ \sin^2\psi(1 - \sin^2\theta\sin^2\varphi) \cdot |A_{\parallel}(t)|^2 \\
 &\quad + \sin^2\psi\sin^2\theta \cdot |A_{\perp}(t)|^2 \\
 &+ (1/\sqrt{2})\sin 2\psi\sin^2\theta\sin 2\varphi \cdot \Re(A_0^*(t)A_{\parallel}(t)) \\
 &+ (1/\sqrt{2})\sin 2\psi\sin 2\theta\cos\varphi \cdot \Im(A_0^*(t)A_{\perp}(t)) \\
 &\quad - \sin^2\psi\sin 2\theta\sin\varphi \cdot \Im(A_{\parallel}^*(t)A_{\perp}(t)).
 \end{aligned}$$

Polarization Amplitudes

Upper sign: Time evolution of pure $B_s^0 \rightarrow J/\psi\Phi$ at $t=0$

Lower sign: Time evolution of pure $\bar{B}_s^0 \rightarrow J/\psi\Phi$ at $t=0$

$$|A_0(t)|^2 = |A_0(0)|^2 \left[\mathcal{T}_+ \pm e^{-\bar{\Gamma}t} \sin \phi_s \sin(\Delta M_s t) \right],$$

$$|A_{\parallel}(t)|^2 = |A_{\parallel}(0)|^2 \left[\mathcal{T}_+ \pm e^{-\bar{\Gamma}t} \sin \phi_s \sin(\Delta M_s t) \right],$$

$$|A_{\perp}(t)|^2 = |A_{\perp}(0)|^2 \left[\mathcal{T}_- \mp e^{-\bar{\Gamma}t} \sin \phi_s \sin(\Delta M_s t) \right],$$

where

$$\mathcal{T}_{\pm} = (1/2) \left[(1 \pm \cos \phi_s) e^{-\Gamma L t} + (1 \mp \cos \phi_s) e^{-\Gamma H t} \right].$$

$$\Re(A_0^*(t)A_{\parallel}(t)) = |A_0(0)||A_{\parallel}(0)| \cos(\delta_2 - \delta_1) \left[\mathcal{T}_+ \pm e^{-\bar{\Gamma}t} \sin \phi_s \sin(\Delta M_s t) \right],$$

$$\Im(A_0^*(t)A_{\perp}(t)) = |A_0(0)||A_{\perp}(0)| \left[e^{-\bar{\Gamma}t} (\pm \sin \delta_2 \cos(\Delta M_s t) \mp \cos \delta_2 \sin(\Delta M_s t) \cos \phi_s) - (1/2) (e^{-\Gamma H t} - e^{-\Gamma L t}) \sin \phi_s \cos \delta_2 \right],$$

$$\Im(A_{\parallel}^*(t)A_{\perp}(t)) = |A_{\parallel}(0)||A_{\perp}(0)| \left[e^{-\bar{\Gamma}t} (\pm \sin \delta_1 \cos(\Delta M_s t) \mp \cos \delta_1 \sin(\Delta M_s t) \cos \phi_s) - (1/2) (e^{-\Gamma H t} - e^{-\Gamma L t}) \sin \phi_s \cos \delta_1 \right],$$

Polarization Amplitudes (no Flavor Tagging)

Assuming equal production rate of B_s^0 and \bar{B}_s^0
 Opposite terms vanish, but still sensitive to φ_s

$$\begin{aligned}
 |A_0(t)|^2 &= |A_0(0)|^2 \left[\mathcal{T}_+ \right], \\
 |A_{\parallel}(t)|^2 &= |A_{\parallel}(0)|^2 \left[\mathcal{T}_+ \right], \\
 |A_{\perp}(t)|^2 &= |A_{\perp}(0)|^2 \left[\mathcal{T}_- \right],
 \end{aligned}$$

where

$$\mathcal{T}_{\pm} = (1/2) \left[(1 \pm \cos \phi_s) e^{-\Gamma L t} + (1 \mp \cos \phi_s) e^{-\Gamma H t} \right].$$

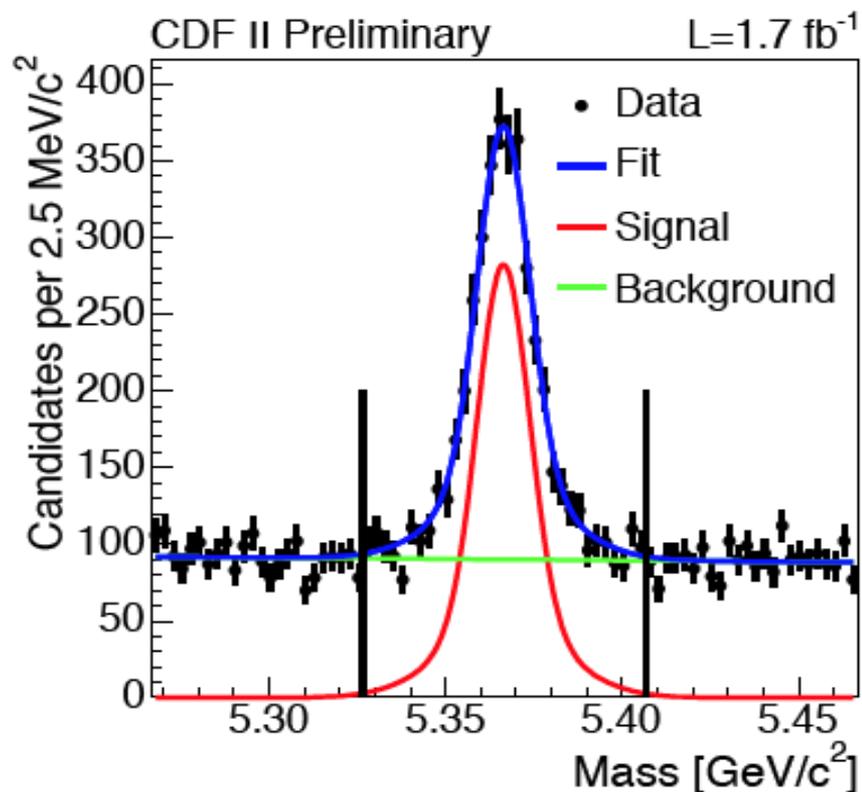
$$\Re(A_0^*(t) A_{\parallel}(t)) = |A_0(0)| |A_{\parallel}(0)| \cos(\delta_2 - \delta_1) \left[\mathcal{T}_+ \right],$$

$$\Im(A_0^*(t) A_{\perp}(t)) = |A_0(0)| |A_{\perp}(0)| \left[e^{-\bar{\Gamma} t} - (1/2) \left(e^{-\Gamma H t} - e^{-\Gamma L t} \right) \sin \phi_s \cos \delta_2 \right],$$

$$\Im(A_{\parallel}^*(t) A_{\perp}(t)) = |A_{\parallel}(0)| |A_{\perp}(0)| \left[e^{-\bar{\Gamma} t} - (1/2) \left(e^{-\Gamma H t} - e^{-\Gamma L t} \right) \sin \phi_s \cos \delta_1 \right],$$

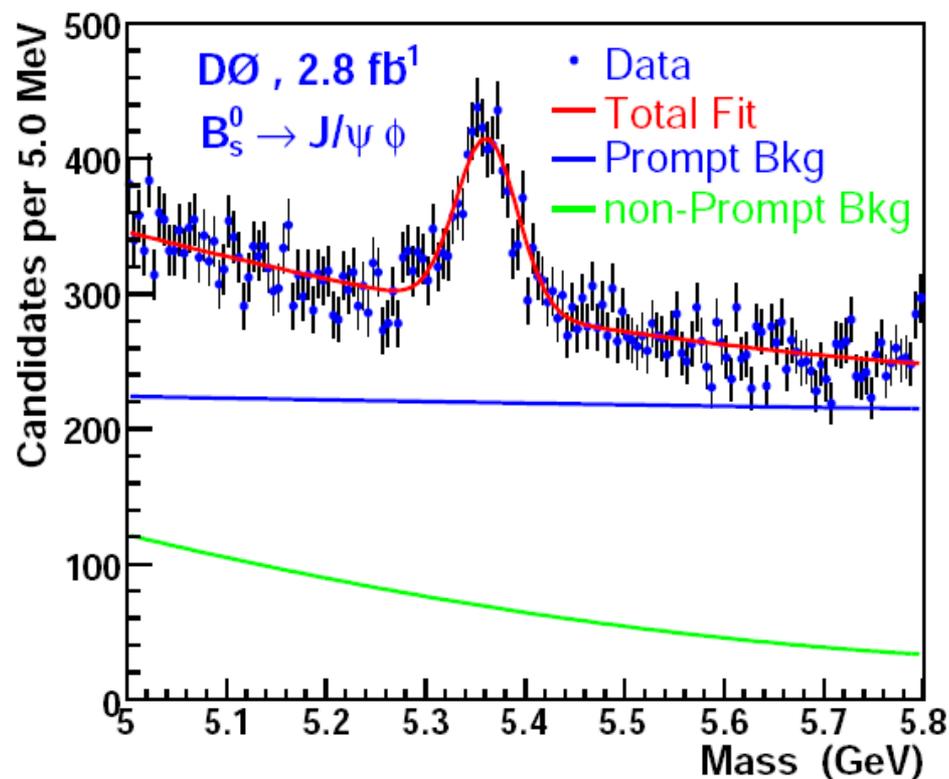
B_s^0 Invariant Mass

CDF



$\sim 2500 B_s^0$ candidates

DØ



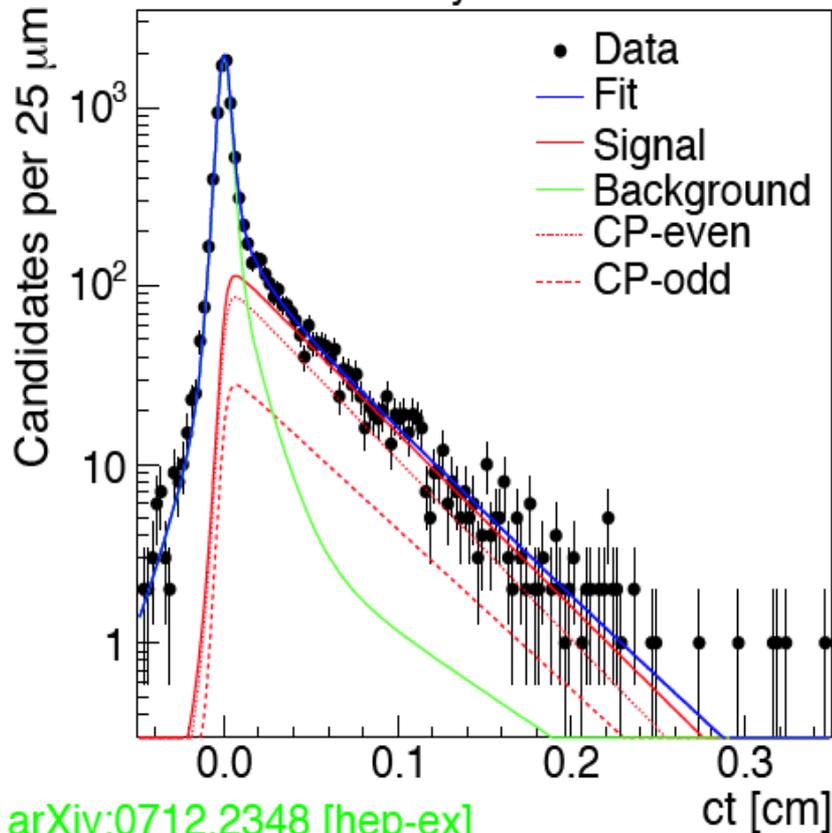
$\sim 2000 B_s^0$ candidates

$\Delta\Gamma_s$ CP Conservation, $\varphi_s = 0$

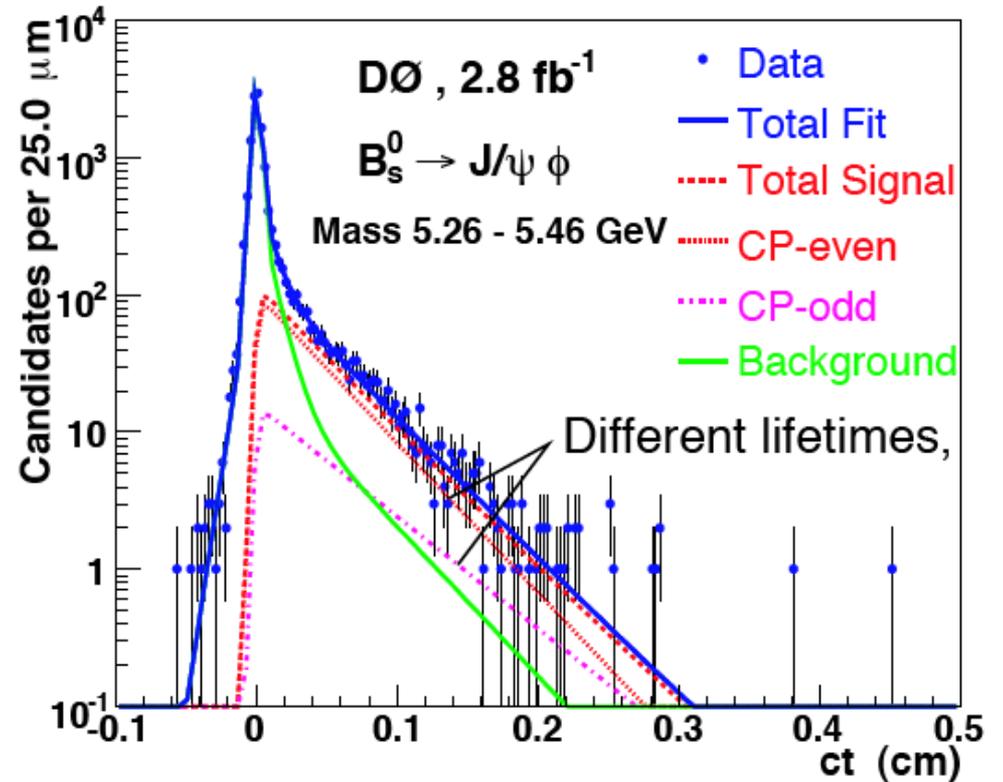
CDF

CDF II Preliminary

$L=1.7 \text{ fb}^{-1}$



DØ



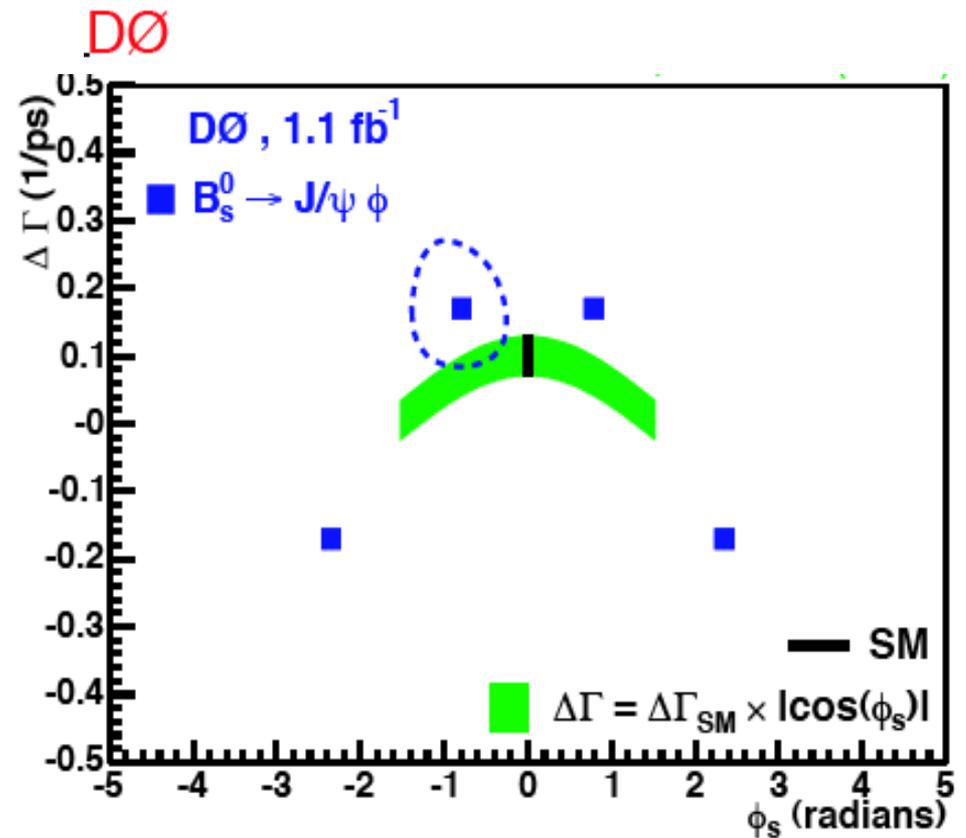
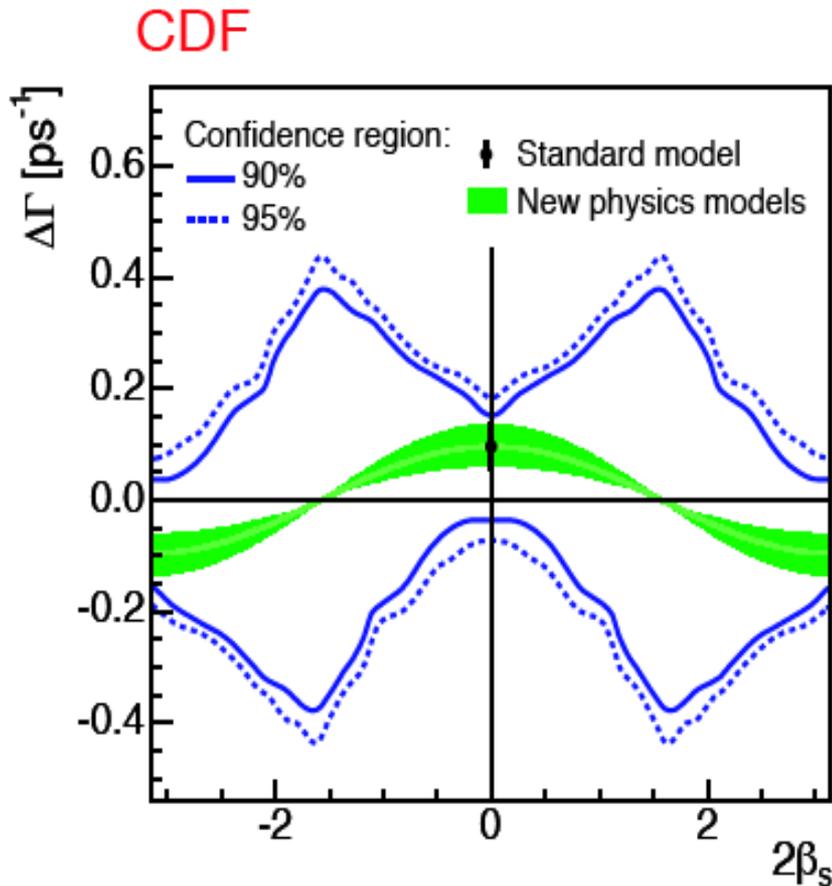
$$\Delta\Gamma_s = 0.076_{-0.063}^{+0.059} \pm 0.006 \text{ ps}^{-1}$$

$$\bar{\tau}_s = 1.52 \pm 0.04 \pm 0.02 \text{ ps}$$

$$\Delta\Gamma_s = 0.14 \pm 0.07 \text{ ps}^{-1} \text{ (Stat + syst)}$$

$$\bar{\tau}_s = 1.53 \pm 0.05 \pm 0.01 \text{ ps}$$

Sensitivity to ϕ_s (no Flavor Tagging)



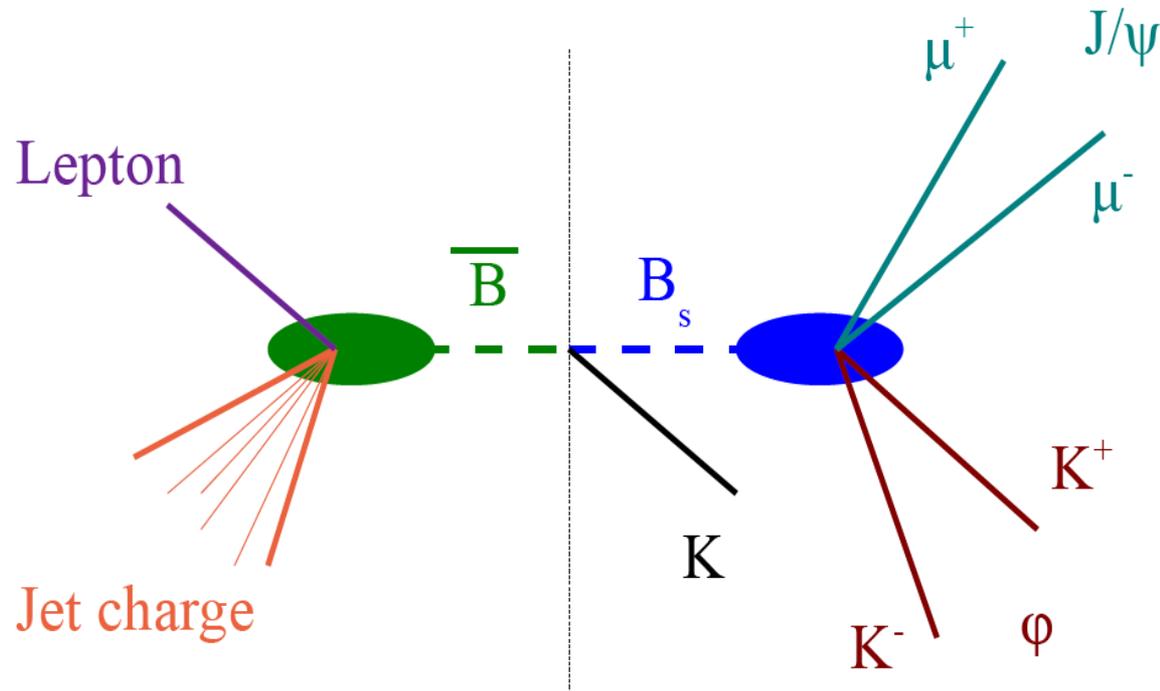
However, four-fold ambiguity reduces to two-fold after applying flavor tagging.

Flavor Tagging

Measurement of B_s^0 or \bar{B}_s^0 flavor at production

Opposite Side

Reconstructed (Same) Side



$$\mathcal{D} \equiv \frac{N_{\text{cor}} - N_{\text{wr}}}{N_{\text{cor}} + N_{\text{wr}}},$$

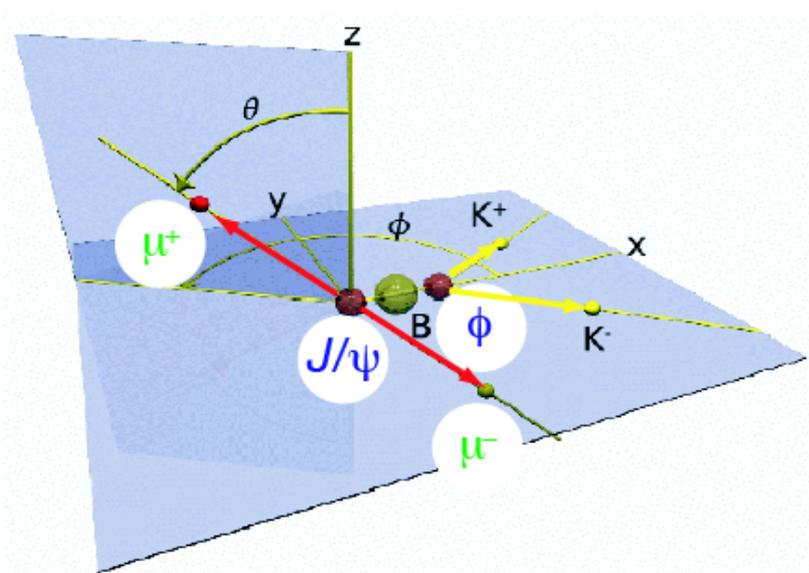
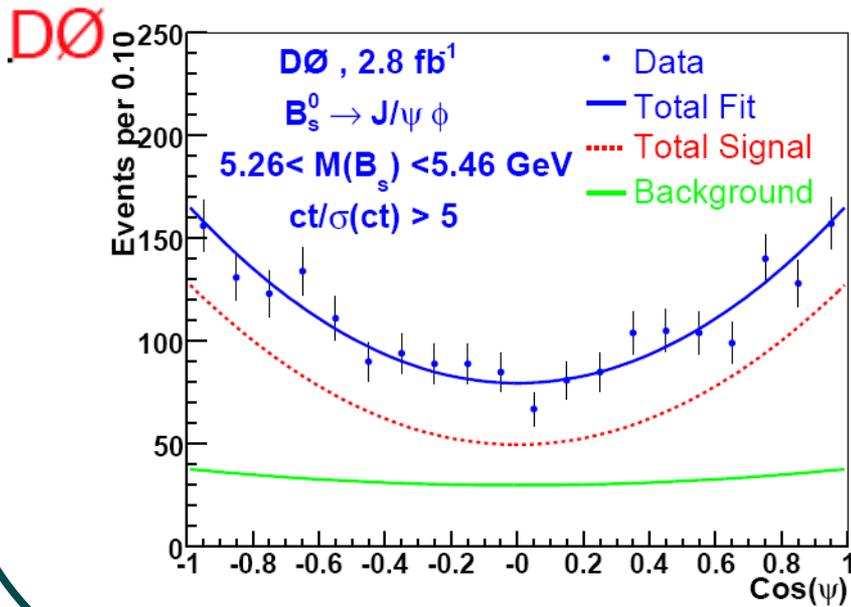
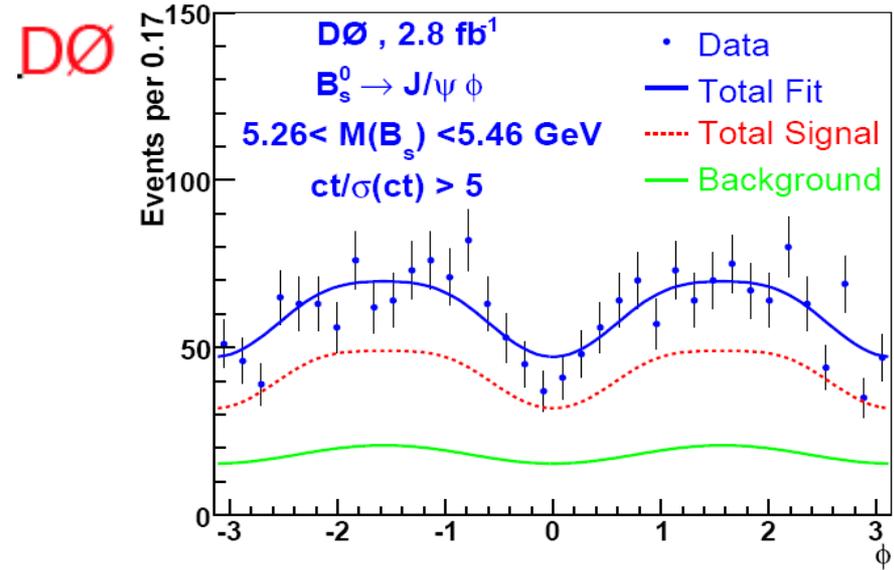
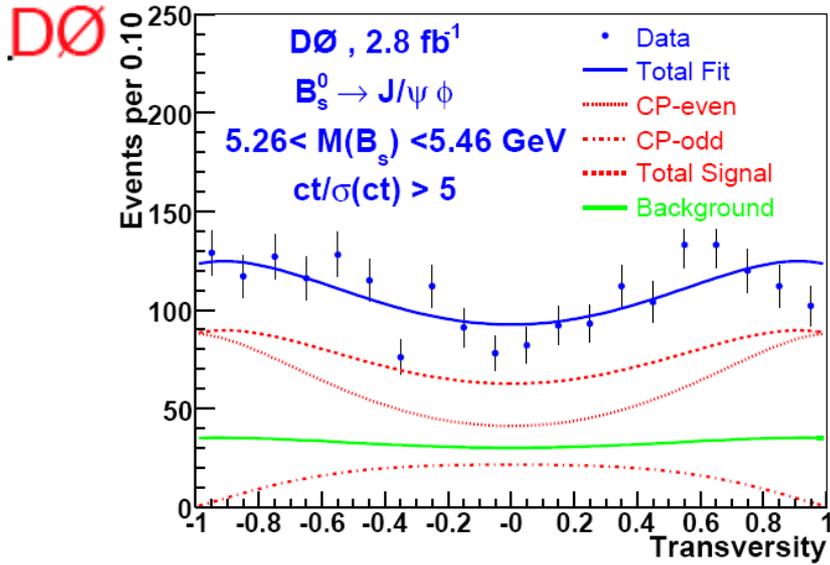
$$\epsilon \equiv \frac{N_{\text{cor}} + N_{\text{wr}}}{N_{\text{tot}}},$$

$$\mathcal{P} \equiv \epsilon \mathcal{D}^2.$$

b quarks produced in pairs

$\epsilon \mathcal{D}^2$ for $B_s^0 \rightarrow J/\psi \phi$ is 4 - 5 %

Angular Fit Projections



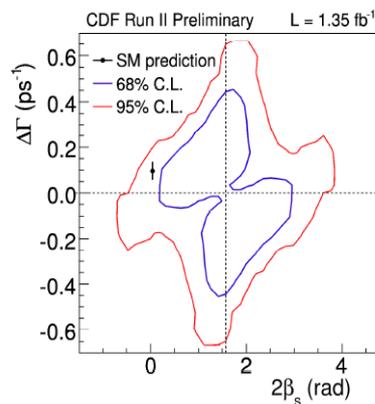
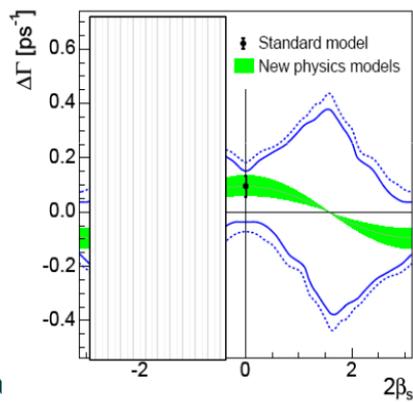


Results with Flavor Tagging

$2\beta_s - \Delta\Gamma$ Confidence Region

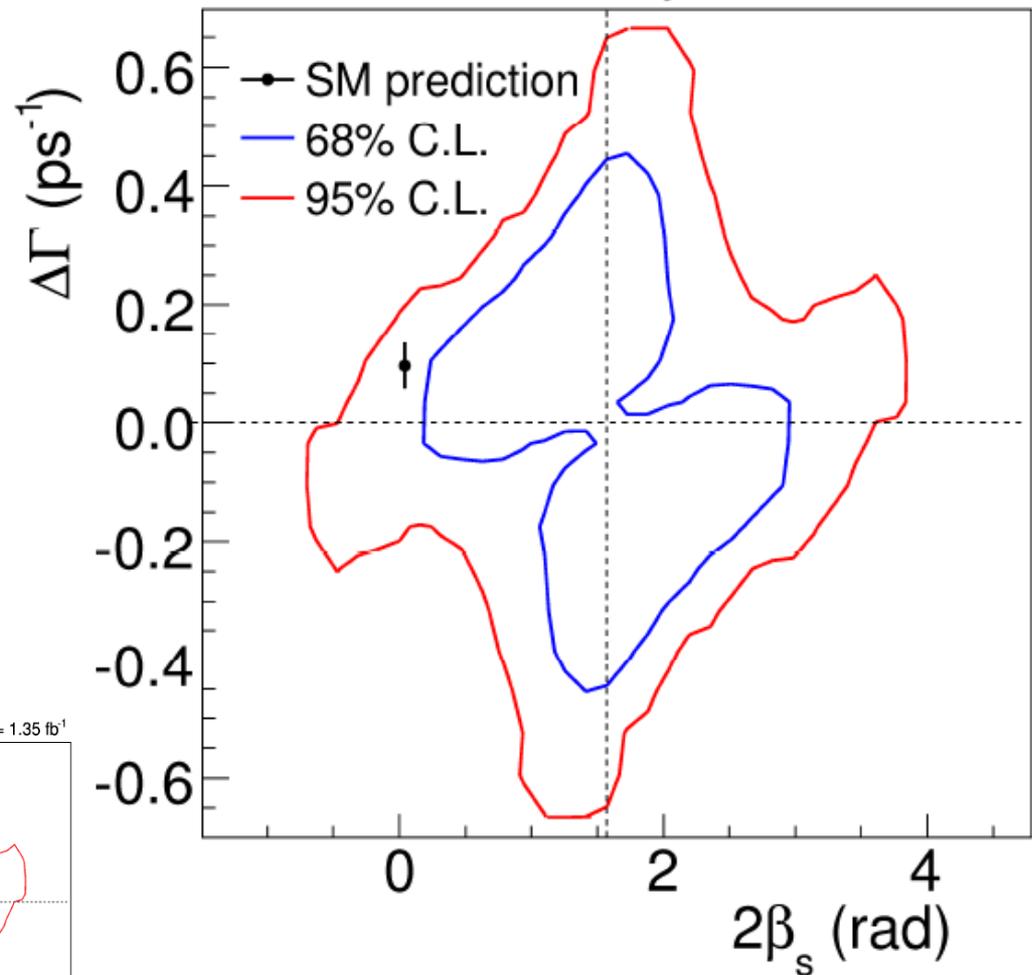
Probability of fluctuation from SM to observation is 15% (1.5σ)

$2\beta_s$ phase space is halved.



CDF Run II Preliminary

$L = 1.35 \text{ fb}^{-1}$

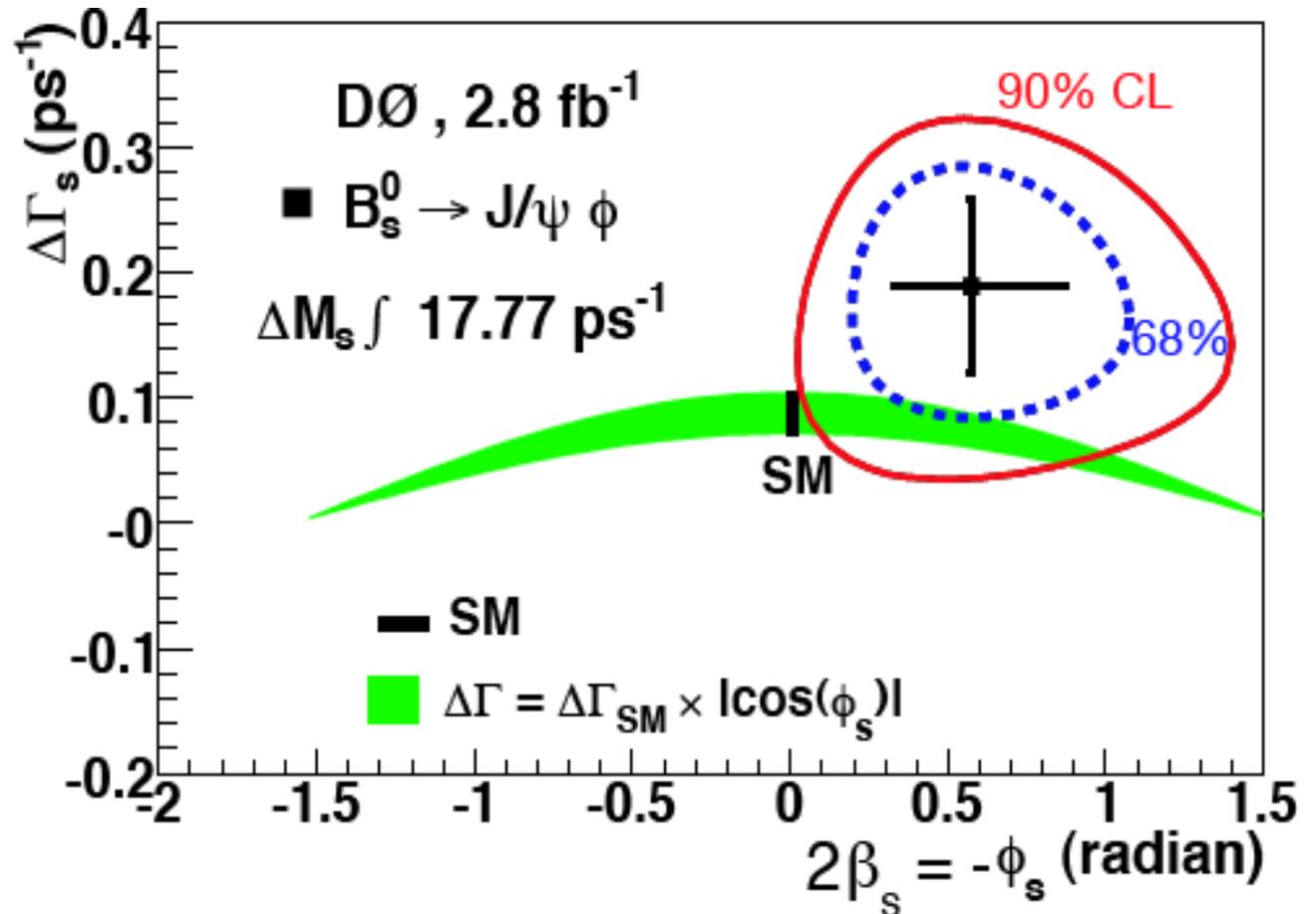


Results with Flavor Tagging



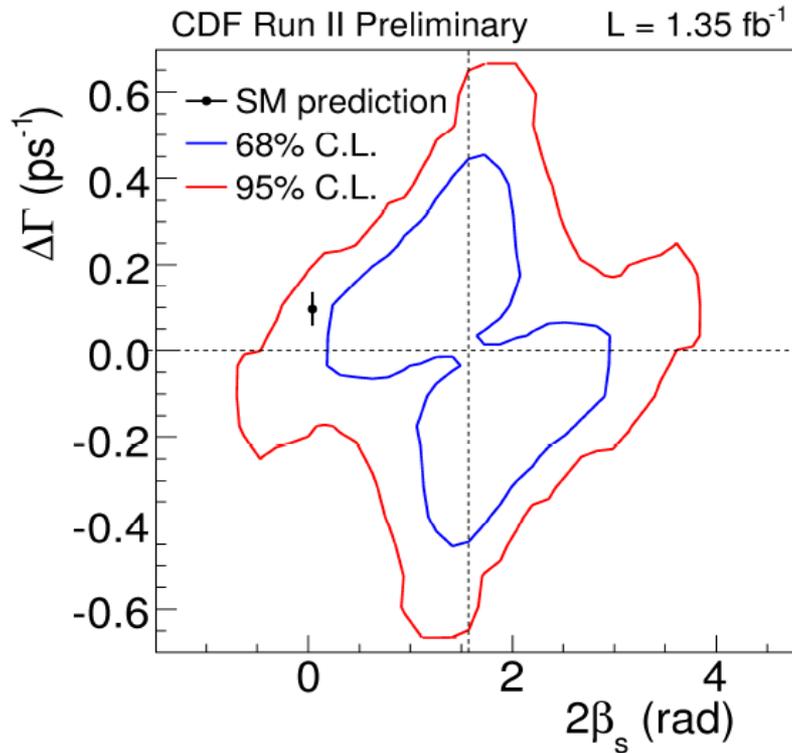
$2\beta_s - \Delta\Gamma$
Confidence Region

Probability of fluctuation
from SM to observation
6.6% (1.8σ)

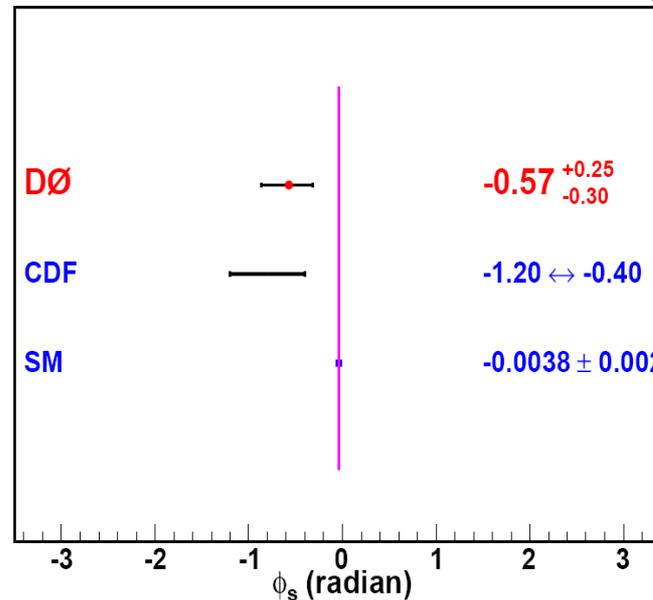
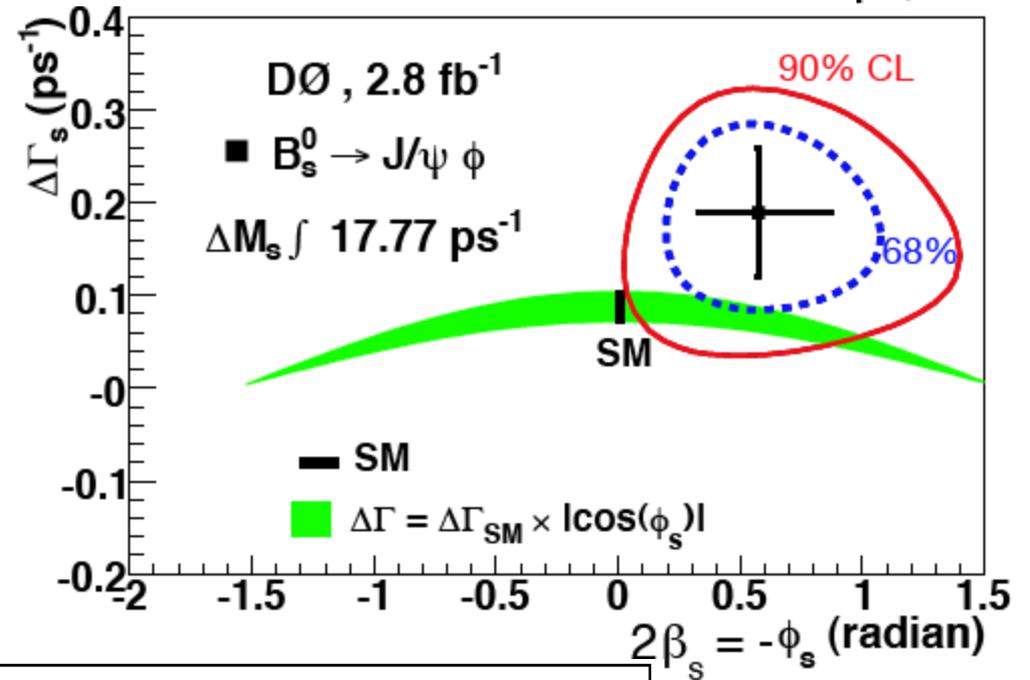


Results with Flavor Tagging

CDF



DØ



Summary

CP Violation is an excellent place to search for new physics beyond the Standard Model.

A large φ_s would be a clear indicator of new physics.

CP studies in the B_s^0 system at CDF and DØ possibly already providing hints of new physics.

A lot of data has been collected and waiting to be analyzed.

Increased datasets will shed more light on the status of CP Violation in the B_s^0 system in the near future.

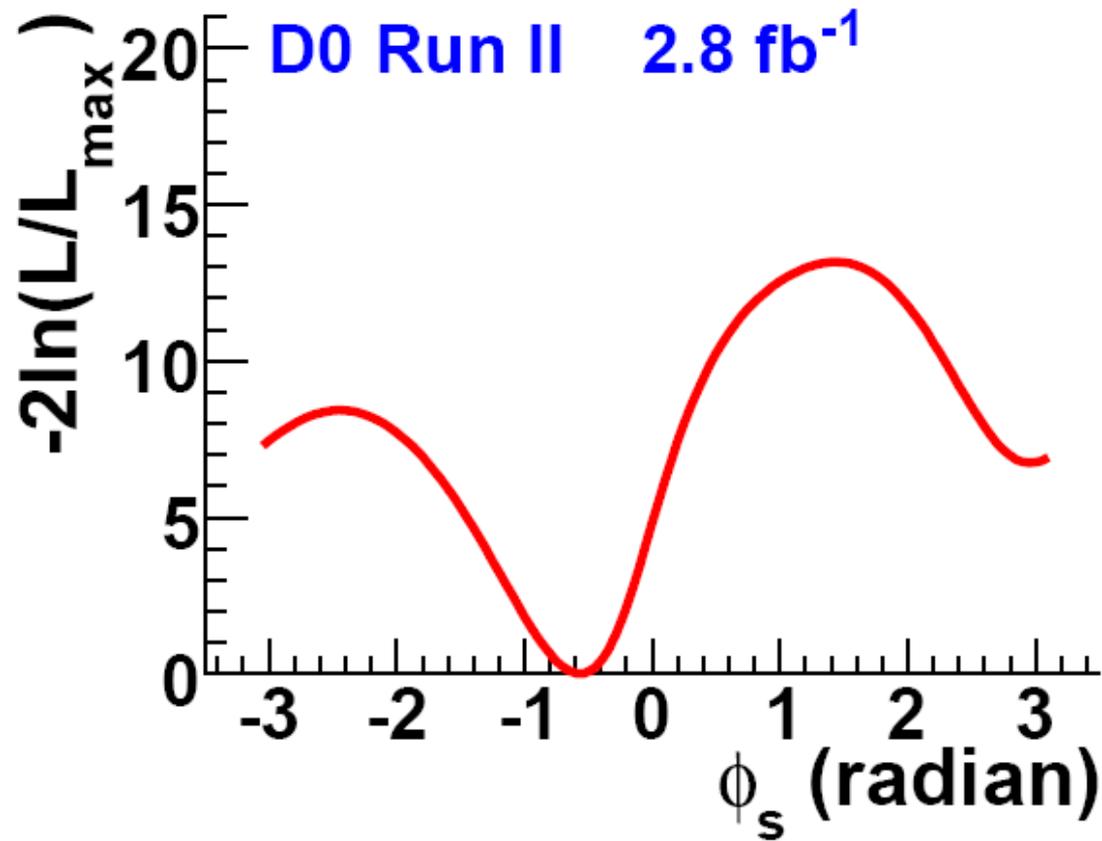
Backup Slides

Systematic Uncertainties

Source	$\bar{\tau}_s$ (ps)	$\Delta\Gamma_s$ (ps ⁻¹)
Acceptance	± 0.003	± 0.003
Signal mass model	-0.01	$+0.006$
Flavor purity estimate	± 0.001	± 0.001
Background model	$+0.003$	$+0.02$
ΔM_s input	± 0.01	± 0.001
Total	± 0.01	$+0.02, -0.01$

Source	$ A_{\perp}(0) $	$ A_0(0) ^2 - A_{\parallel}(0) ^2$	ϕ_s
Acceptance	± 0.005	± 0.03	± 0.005
Signal mass model	-0.003	-0.001	-0.006
Flavor purity estimate	± 0.001	± 0.001	± 0.01
Background model	-0.02	-0.01	$+0.02$
ΔM_s input	± 0.001	± 0.001	$+0.06, -0.01$
Total	$+0.01, -0.02$	± 0.03	$+0.07, -0.02$

Likelihood Profile



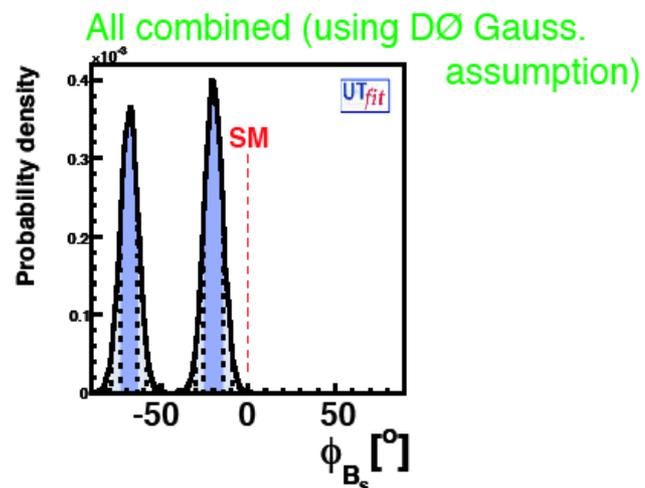
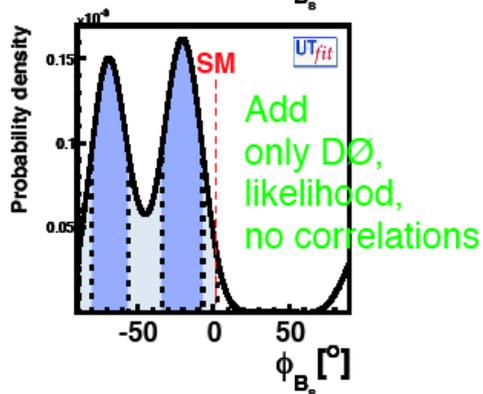
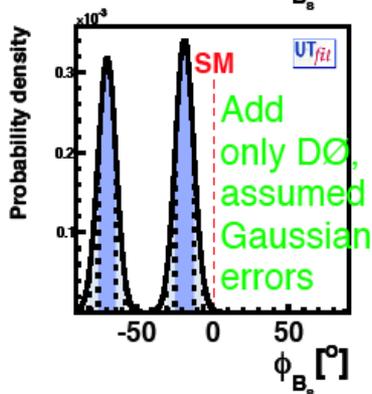
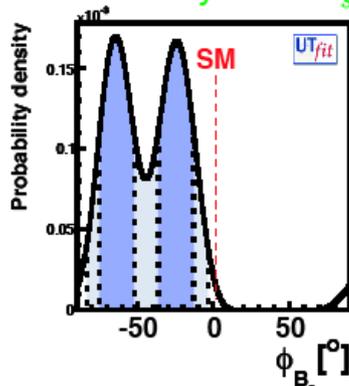
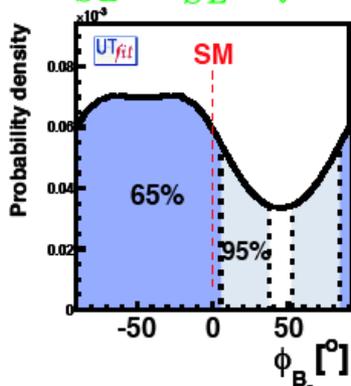
Likelihood profile of ϕ_s .

Combination

UTfit group, arXiv:0803.0659:

$\Delta m_s, A_{SL}^s, A_{SL}^{\mu\mu}, \tau_{fs}$

Add only CDF $B_s^0 \rightarrow J/\psi\phi$



> 3σ deviation from SM

- Intriguing, results from D0, CDF with more data needed

New HFAG Result

1-sigma contours ($\Delta(\log L) = 0.5$)

