

## Fermion polarisation probes MSSM w/o CP conservation at LC.

- ◇ Introduction.
- ◇  $\tau$  polarisation in  $\tilde{\tau}$  decays at  $e^+e^-$  colliders: CP conserving case.
- ◇  $\tau(t)$  polarisation in  $\tilde{\tau}(\tilde{t})$  decays at the  $e^+e^-$  colliders: CP violating case.
- ◇ CP mixing induced in the Higgs sector due to CP violation in soft SUSY breaking parameters.
- ◇ Probing CP violation in the Higgs sector through  $\tau$  polarisation at  $\gamma\gamma$  colliders.
- ◇ Conclusions.

Some of the refs:

S. Kraml, D. Miller, M. Krawczyck and R. Godbole Contribution to the LHC/LC studies, [arXiv:hep-ph/0404024].

M. Guchait, D.P. Roy and R.G. (Manuscript in preparation)

S. Kraml, T. Gadosijk, R.G., JHEP **0409**, 051 (2004) [arXiv:hep-ph/0405167].

S. Kraml, R.K. Singh and R.G. [arXiv:hep-ph/0409199], Talk given at International Conference on Linear Colliders (LCWS 04), Paris, France, 19-24 Apr 2004, (Manuscript in preparation.)

## Introduction

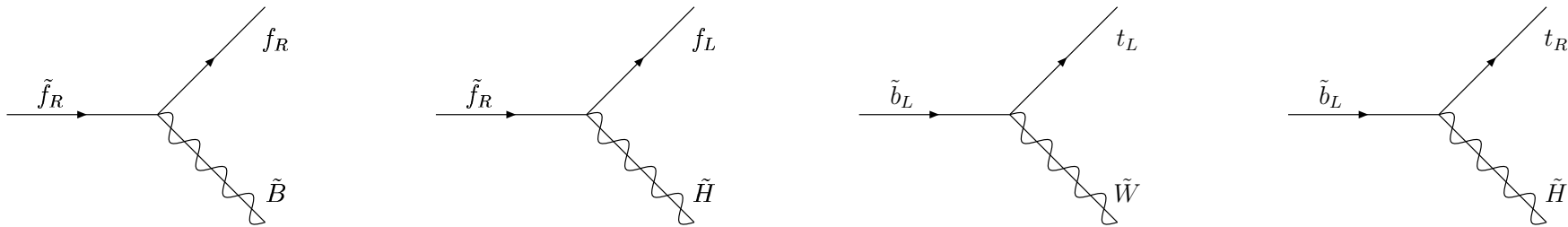
### Why study $\tau / (t)$ polarisation?

- Large mass of the third generation fermions means large coupling to a Higgs. Third generation fermions useful to probe Higgs physics.
- $f\bar{f}$  pair produced via gauge interactions will have opposite handedness, produced via Yukawa interactions will have the same handedness. A correlation between polarisation of  $f$  and  $\bar{f}$  can be a probe of Higgs contribution.
- $\tau / (t)$  polarisation probes chirality structure of  $\tau(t)$  production process, may it be Higgs physics, new physics : eg.  $H^\pm \rightarrow \tau\nu_\tau$  decay Vs  $W \rightarrow \tau\nu_\tau$ .  $\tau$ 's produced in the former(latter) are right(left) handed, modulo  $m_\tau$  effects, CP violation in the Higgs sector will reflect in the decay  $\tau / (t)$  polarisation.

- Large mass of the top  $\Rightarrow t$  decays before hadronisation. The decay  $l$  can retain the memory of the  $t$  polarisation.
  - $\tau$  has hadronic decay modes. The energy distribution of the  $\pi$  produced in the decay,  $\tau \rightarrow \nu_\tau \pi$  as well as those in  $\tau \rightarrow \rho \nu_\tau, \tau \rightarrow A_1 \nu_\tau$  depends on the handedness of the  $\tau$ . Thus  $\tau$  polarisation can be determined using decay  $\pi$  energy distribution. K. Hagiwara, A.D. Martin and D. Zeppenfeld, PLB **235** 198 (1990), B.K.Bullock, K.Hagiwara and A.D.Martin, PRL **67**, 3055 (1991), NPB **395**, 499 (1993).
  - If  $m_h < 2m_t$  then  $\tau$  is the only fermion whose polarisation can be determined and into which a Higgs decays may be somewhat substantial.
- ◇ Thus  $t$  and  $\tau$  polarisation are a very good probe for new and Higgs physics.

$\tau(t)$  produced in stau/stop decay. M. Nojiri, PRD 51 (1995) 6281 [hep-ph/9412374]

$$f = t/\tau$$



- In MSSM mass eigenstates  $\tilde{f}$ -sleptons/squarks ( $\tilde{f}_1, \tilde{f}_2$ ), mixtures of  $\tilde{f}_L$  and  $\tilde{f}_R$ ,  $f = t, \tau$ .
- Mixing affects Gauge couplings of  $\tilde{f}_i, i = 1, 2$  and hence the production rates.
- The  $\tilde{\chi}_j^\pm, j = 1, 2, \tilde{\chi}_j^0, j = 1, 4$  are mixtures of Higgsinos and Gauginos.
- Couplings of sfermions with Higgsinos flip chirality whereas those with gauginos do not.
- Net helicity of produced  $f$  in the decay  $\tilde{f}_i \rightarrow \tilde{\chi}_j^0 f$  AND  $\tilde{f}_i \rightarrow \tilde{\chi}_j^\pm f'$  depends on the  $L-R$  mixing in the sfermion sector and gaugino-higgsino mixing.

- A lot of analyses of  $\tau$  polarisation and hence of the MSSM parameter determination at LC exist. Use the  $\tau \rightarrow \rho/A\nu_\tau$  (multiprong) mode.

- Our New work: (M. Guchait, D.P. Roy and R.G.) Developed a new variable for  $\tau$  polarisation analysis.

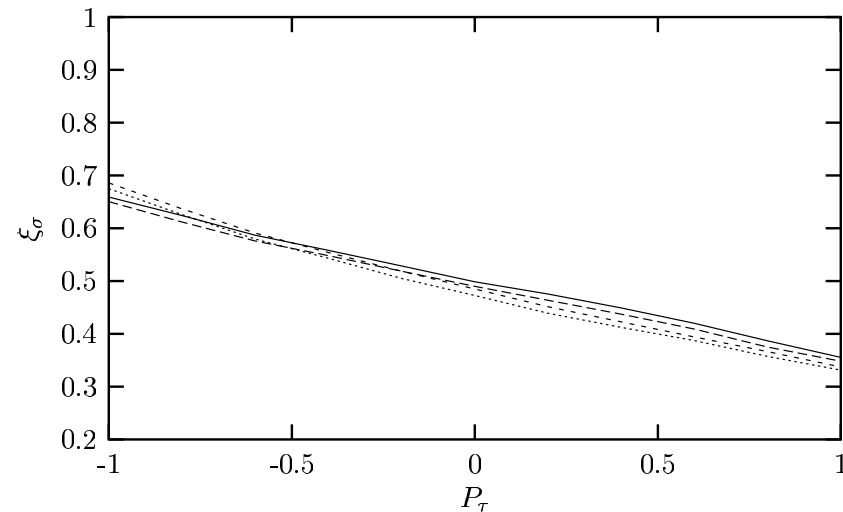
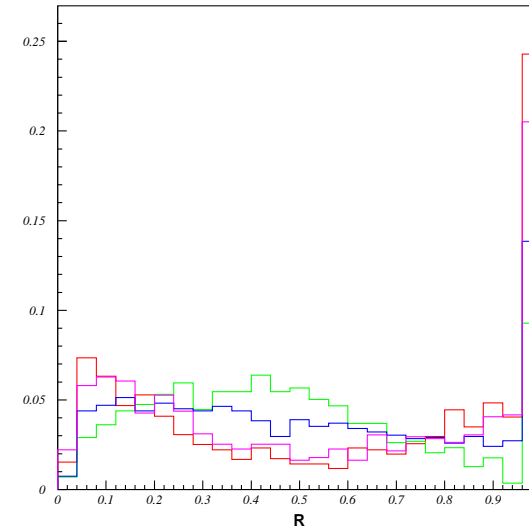
- Look at  $R = \frac{E_\pi}{E_{jet}}$ . and study

$$\xi = \frac{\#(0.2 < R < 0.8)}{\#total}$$

- $\xi$  a good discriminator of  $\tau$  polarisation, hence of SUSY models.

Top :Distribution in R. Different curves for values of polarisation  $p_\tau = -1, -0.5, 0.5, 1.0$ . Bottom:  $\xi$  as a function of  $P_\tau$ , band: theoretical uncertainty, obtained using Tauola.

- Detailed application of this variable to MSSM parameter analysis in progress.



- Polarisation of  $f(f')$  produced in  $\tilde{f}_i \rightarrow f\tilde{\chi}_j^0$ ,  $\tilde{f}_i \rightarrow \tilde{\chi}_j^\pm f'$ , depends on  $L-R$  mixing and Guagino/higgsino content
- ◇ CP-violating phases of  $A_f, \mu, M_i$  in MSSM affect  $L-R$  mixing and Guagino/higgsino content
- Polarisation of  $f, f'$  can carry information on CPV phases as well.
- ◇ NOTE: polarisation itself is a CP-even variable
- Our study: Kraml, T. Gadosijk, R.G JHEP 0409, 051 (2004) How well does the  $P_f$  probe the CPV phases?

- An example for  $\tilde{\tau}_1 \rightarrow \tau \tilde{\chi}_1^0$

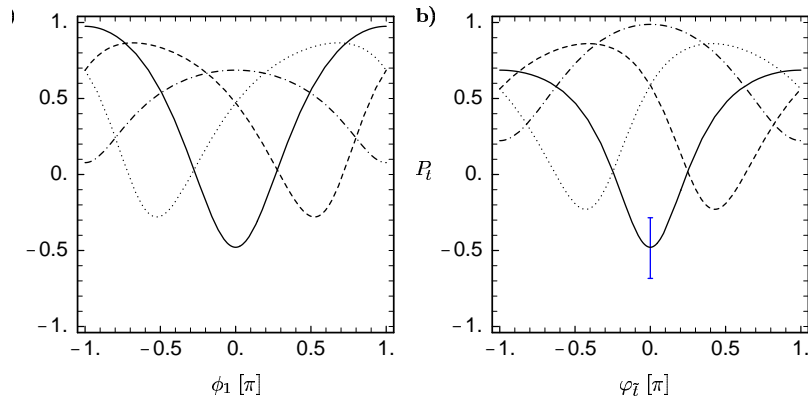
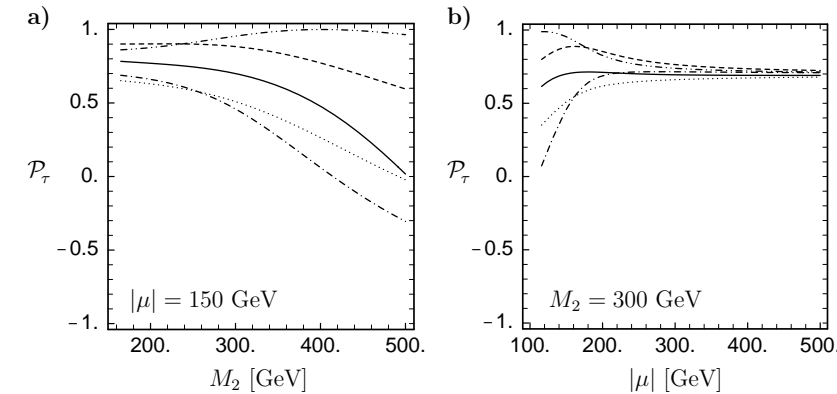
$$P_f = \frac{Br(\tilde{f}_1 \rightarrow \tilde{\chi}_1^0 f_R) - Br(\tilde{f}_1 \rightarrow \tilde{\chi}_1^0 f_L)}{Br(\tilde{f}_1 \rightarrow \tilde{\chi}_1^0 f_R) + Br(\tilde{f}_1 \rightarrow \tilde{\chi}_1^0 f_L)}$$

- Take  $A_f, M_2, M_1, \mu$  complex. Safe to choose  $\mu$  real  $\Leftarrow$  (EDM CONSTRAINTS).
- ◇  $P_f$  sensitive to CPV and nonzero even if ONLY one phase (either in sfermion sector or gaugino sector) is nonzero.  $m_{\tilde{f}}$  not relevant for predictions of  $P_f$ .
- Effects large for larger Yukawa Couplings and enhanced if  $\tilde{\chi}_1^0$  is a gaugino-higgsino mixed state.
- $P_t$  has significant dependence on the phases even when  $\mu \simeq M_2$
- ◇  $P_f$  can be used to extract information on phases ONLY in conjunction with other observables which will give info. on  $\mu, M_2, M_1$  etc.
- ◇ Polarisation information need be included in a global analysis of MSSM parameter determination at ILC.



## Some Results:

Average polarisation for  $\tau$  for  $\theta_{\tilde{\tau}} = 130^\circ$  and  $\tan\beta = 10$ : in a) as a function of  $M_2$  for  $|\mu| = 150$  GeV, in b) as a function of  $|\mu|$  for  $M_2 = 300$  GeV. The full, dashed, dotted, dash-dotted, and dash-dot-dotted lines are for  $(\phi_1, \varphi_{\tilde{\tau}}) = (0, 0)$ ,  $(0, \frac{\pi}{2})$ ,  $(\frac{\pi}{2}, 0)$ ,  $(\frac{\pi}{2}, \frac{\pi}{2})$ , and  $(\frac{\pi}{2}, -\frac{\pi}{2})$ .



Average polarisation of top for  $\theta_{\tilde{t}} = 130^\circ$ , and  $\tan\beta = 10$ : in a) as a function of  $\phi_1$  for  $M_2 = 225$  GeV and  $|\mu| = 200$  GeV; in b) as a function of  $\varphi_{\tilde{\tau}}$  for  $|\mu| = 200$  GeV and  $M_2$  adjusted such that  $m_{\tilde{\chi}_1^0} = 100$  GeV. The full, dashed, dotted, and dash-dotted lines are for  $\varphi_{\tilde{\tau}} (\phi_1) = 0, \frac{\pi}{2}, -\frac{\pi}{2}, \pi$  in a (b).

MSSM  $\mathcal{CP}$  phases  $\Rightarrow$   $\mathcal{CP}$  in the Higgs sector:

$CP$  conserving MSSM Three Neutral Higgses  $\begin{matrix} h, H \\ CP\text{-even} \end{matrix}$   $\begin{matrix} A \\ CP\text{-odd} \end{matrix}$

$CP$  violation :  $\begin{matrix} \phi_1, \phi_2, \phi_3 \\ \text{no fixed } CP \text{ property} \end{matrix}$

$$m_{\phi_1} < m_{\phi_2} < m_{\phi_3}$$

The  $h, H, A$  now all mix and share the couplings with vector boson pair  $VV$ . Will affect production rates.

$$g^{VV\phi_1} < g^{VVH_{SM}} \Rightarrow \sigma(e^+e^- \rightarrow Z^* \rightarrow Z\phi) < \sigma(e^+e^- \rightarrow Z^* \rightarrow ZH_{SM})$$

May escape detection at  $e^+e^-$  collider, but can still be produced at a  $\gamma\gamma$  collider.

Predictions in terms of SUSY  $\mathcal{CP}$  phases in the MSSM for this mixing.

Tools

## CPSuperH

Lee, Pilaftsis, Carena, Choi, Drees, Ellis &amp; Wagner

<http://theory.ph.man.ac.uk/~jslee/CPsuperH.html>

## FeynHiggs

Hahn, Heinemeyer, Hollik &amp; Weiglein

<http://www.feynhiggs.de>low energy parameters  $\longrightarrow$  masses, BR's, couplings...RG improved effective potential  
for masses & couplings

leading log approx for one-loop

leading log approx for  $O(\alpha_s\alpha_t, \alpha_t^2)$ ,  
but full phase dependenceFeynman–diagrammatic approach  
for masses & couplings

full one-loop

full  $O(\alpha_s\alpha_t, \alpha_t^2)$  but  
approx phase dependence $O(\alpha_s\alpha_t)$  has  $(\alpha_s \tan \beta)^n$  resummation and full complex phase  
dependence

CPX Scenario [Carena, Ellis, Pilaftsis & Wagner, Phys. Lett. **B495** (2000) 155]

“designed to showcase the effects of CP violation in the MSSM”

$$M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = M_{\tilde{L}_3} = M_{\tilde{E}_3} = M_{\text{SuSy}}$$

$$\mu = 4M_{\text{SuSy}}, \quad |A_{t,b,\tau}| = 2M_{\text{SuSy}}, \quad |M_3| = 1\text{TeV}, \quad |M_1| = |M_2| = 200\text{GeV}$$

Allow the following parameters to vary:

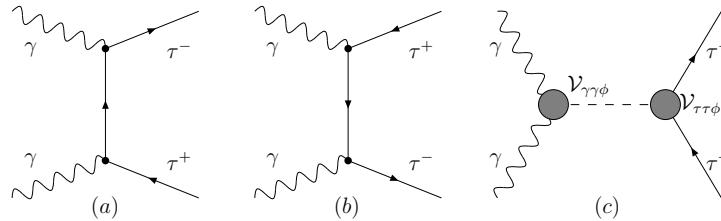
$\tan \beta,$	$M_{H^\pm},$	$M_{\text{SuSy}},$
$\{\Phi_{A_t}, \Phi_{A_b}, \Phi_{A_\tau}\},$	$\Phi_3,$	$\Phi_\mu$

What have we done?

Kraml, R.Singh, R.G, [arXiv:hep-ph/0409199]

$\tau$  pair production as a probe of the Higgs contribution and  $CP$  structure of its coupling. Studied

$$\gamma\gamma \rightarrow \phi \rightarrow \tau^+\tau^-$$



The Higgs contribution treated in a model independent way.

$$\mathcal{V}_{\tau^+\tau^-\phi} = -ie \frac{m_\tau}{M_W} (S_\tau + i\gamma^5 P_\tau),$$

$$\mathcal{V}_{\gamma\gamma\phi} = \frac{-i\sqrt{s}\alpha}{4\pi} \left[ S_\gamma(s) \left( \epsilon_1 \cdot \epsilon_2 - \frac{2}{s} (\epsilon_1 \cdot k_2)(\epsilon_2 \cdot k_1) \right) - P_\gamma(s) \frac{2}{s} \epsilon_{\mu\nu\alpha\beta} \epsilon_1^\mu \epsilon_2^\nu k_1^\alpha k_2^\beta \right].$$

$\{S_\tau, P_\tau, S_\gamma, P_\gamma\}$  depend upon the MSSM parameters.

Combinations of form-factors that appear in the helicity amplitude.

Combinations	Aliases	CP-property	Combinations	Aliases	CP-property
$S_\tau \Re(S_\gamma)$	$x_1$	even	$S_\tau \Re(P_\gamma)$	$y_1$	odd
$S_\tau \Im(S_\gamma)$	$x_2$	even	$S_\tau \Im(P_\gamma)$	$y_2$	odd
$P_\tau \Re(P_\gamma)$	$x_3$	even	$P_\tau \Re(S_\gamma)$	$y_3$	odd
$P_\tau \Im(P_\gamma)$	$x_4$	even	$P_\tau \Im(S_\gamma)$	$y_4$	odd

QED background :  $P$ ,  $CP$  and chirality conserving.

Higgs exchange diagram violates these symmetries,

$\{x_i, y_j\} \neq 0 \Rightarrow$ . Chirality flipping interaction,  $\Rightarrow$   $\tau$ -polarisation affected.

Polarisation of  $\tau$ :

$$P_{\tau}^{IJ} = \frac{N_{+}^{IJ} - N_{-}^{IJ}}{N_{+}^{IJ} + N_{-}^{IJ}}$$

$I, J = +, -, U$  Polarisation of parent  $e^+/e^-$  beam.

$N_{+}^{IJ} = \#$  of  $\tau_R$ ,  $N_{-}^{IJ} = \#$  of  $\tau_L$ .

◇  $P_{\tau}^U = 0$  : for QED contribution.

◇  $P_{\tau}^U = 0$  : even with Higgs contribution, if  $y_j = 0$ ; i.e. Higgs is a  $CP$  eigenstate.

□  $P_{\tau}^U \propto m_{\tau} \Rightarrow$  small **even** with CPV.

□  $P_{\tau}^U$  a 'poor' but 'pure' probe of CPV in the Higgs sector.

Move to polarised photons :

- $P_\tau^{++}, P_\tau^{--}$  : finite for QED diagrams alone.

$P$  invariance of QED  $\Rightarrow P_\tau^{++} = -P_\tau^{--} \Rightarrow P_\tau^{++} + P_\tau^{--} = 0$

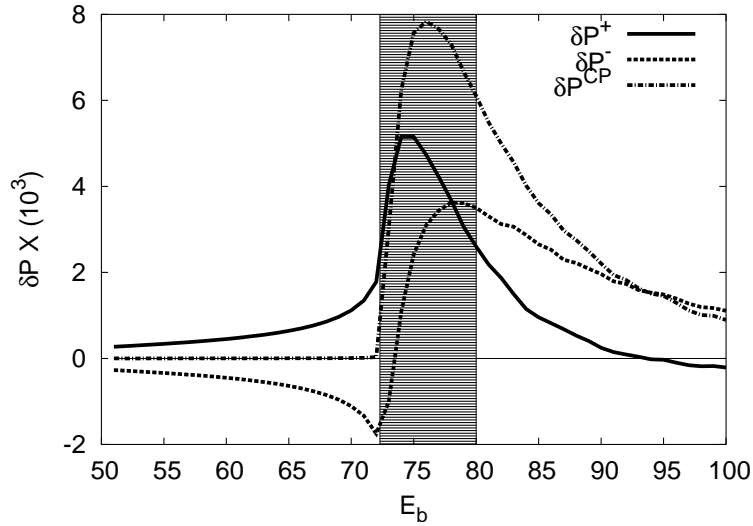
$P_\tau^{++} + P_\tau^{--} \neq 0 \Rightarrow$  signal of  $\not{P}$  In case of  $C$  invariance  $\Rightarrow$  signal of  $CP$  violation.

- $P_\tau^{++}$  : modified by the Higgs contribution.

- $P_\tau^{++} - (P_\tau^{++})^{QED} \neq 0$  even if  $\phi$  is  $CP$  eigenstate,  $\Rightarrow$  probe of chirality flipping amplitude.

Observables	Interaction probed
$P_\tau^U$	CP-violating interaction
$\delta P_\tau^{CP} = P_\tau^{++} + P_\tau^{--}$	CP-violating interaction
$\delta P_\tau^+ = P_\tau^{++} - (P_\tau^{++})^{QED}$	Chirality flipping interaction
$\delta P_\tau^- = P_\tau^{--} - (P_\tau^{--})^{QED}$	Chirality flipping interaction





$\delta P_{\tau}^{+}, \delta P_{\tau}^{CP}$  as a function of  $E_b$

$m_{\phi_1} = 120$  GeV for a chosen CP violating MSSM point.

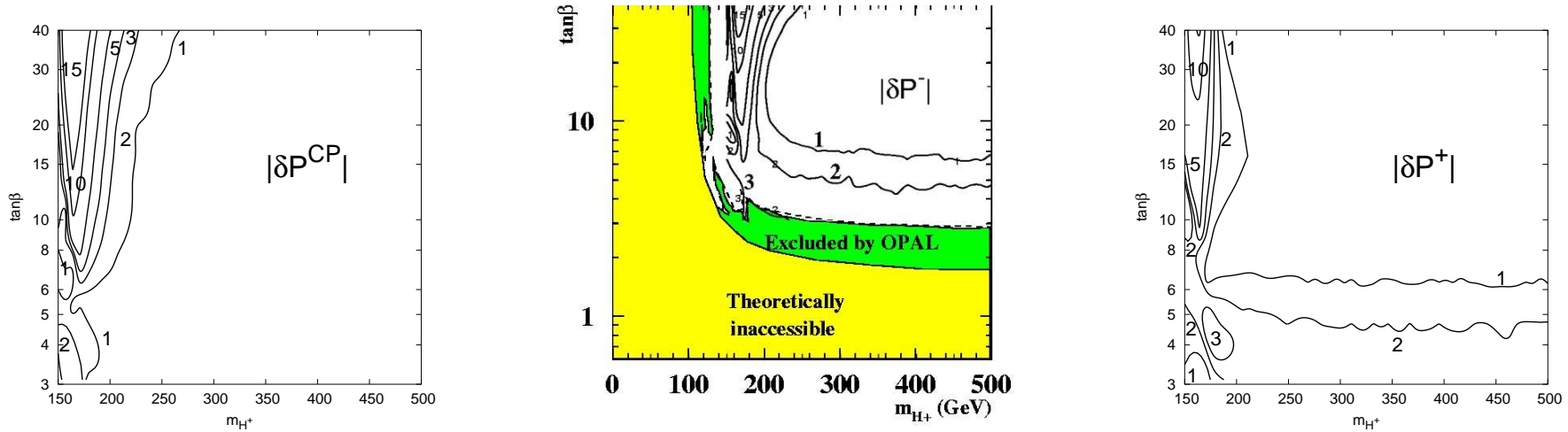
Ideal back-scattered photons,  $x_c = 4.8$ .

$\delta P_{\tau}^{+}, \delta P_{\tau}^{CP}$  peak for  $E_b \approx (m_h/2)/f$  where  $0.75 < f < 0.83$ . This is the grey region.

Beam energy can be chosen to maximise the sensitivity

Main contribution from off-shell Higgs  $\Rightarrow$  large statistics available.

Contribution from  $\phi_2, \phi_3$  may be important if they are close to  $\phi_1$  in mass.



- Contribution from all the three  $\phi_i$  included. Over the scanned region of the parameter space the degeneracy of the three states not high, i.e.  $\delta m = m_{\phi_i} - m_{\phi_j} \gg \Gamma_{\phi_i} + \Gamma_{\phi_j}$
- $\delta P_{\tau}^{CP}$  is sensitive to some part of region.  $\delta P_{\tau}^-$  is sensitive to most of the region. Only two of the three  $P_{\tau}^X$  linearly independent.
- If sensitivity of  $P_{\tau}^{IJ} \approx 0.001$ , the observables can probe all the OPAL allowed region of the MSSM parameter space.

## Conclusions:

- $\tau/t$  polarisation a very useful probe of different aspects of Higgs physics.
- Fraction of events in the inclusive pion spectrum with  $0.2 < R < 0.8$  where  $R = \frac{E_\pi}{E_{jet}}$ , is correlated nicely with  $\tau$  polarisation. This is a new observable to measure the  $\tau$  polarisation.
- CP violating phases in the MSSM affect the CP-even polarisation of  $\tau(t)$  produced in stau/stop/sbottom decays. Effects larger for the  $t$  due to larger Yukawa coupling. Effects larger when  $\tilde{\chi}_j^0, \tilde{\chi}_i^\pm$  is a mixed gaugino-higgsino state. Determination of CP phases requires combining the polarisation information with knowledge of magnitudes of MSSM parameters from other observables.
- Program of global determination of SUSY parameters at the ILC should include polarisation information.
- Have constructed probes of CP violation and chirality flipping interactions in  $\gamma\gamma \rightarrow \tau^+\tau^-$ .
- ◇  $\delta P_\tau^-$  requires one measurement,  $\delta P_\tau^{CP}$  requires two measurements. Both cover large regions of the OPAL allowed MSSM parameter space in CPX scenario.
- ◇  $\tau$  polarisation as a probe of CP properties of neutral Higgs holds promise